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GENERAL ATOMIC DIVISION OF GENERAL DYNAMICS

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OCD Work Unit # 3111A

A CALCULATIONAL MODEL FOR CONDENSED STATE DIFFUSION CONTROLLED FISSION FRODUCT ABSORPTION **DURING FALLOUT FORMATION**

by

R. F. Korts and J. H. Norman

Performed for Office of Civil Defense Secretary of the Army Department of Defense, 20301 under Contract N0022866C0403 through the U.S. Naval Raciological Defense Laboratory San Francisco, California 94135



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January 10, 1967



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The FORTRAN IV computer code for the UNIVAC 1108 described herein is as it existed on 12/1/66. The code has been in continuous development for 1 year and in its presented form has been applied successfully by General Atomic to the kind of problems discussed later in this report. However, the development and improvement of the code are being continued, so that duplication of results (or even close agreement) between problems run with the code as published and the code as it existed either before or after this time is not necessarily to be expected.

General Atomic has exercised due care in preparation, but does not warrant the merchantability, accuracy, and completeness of the code or of its description contained herein. The complexity of this kind of program precludes any guarantee to that effect. Therefore, any user must make his own determination of the suitability of the code for any specific use, and of the validity of the information produced by use of the code.

GENERAL ATOMIC

DIVISION OF

GENERAL DYNAMICS

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This report has been reviewed in OCD and approved for publication. Approval does not signify that the contents necessarily reflect the views and policies of the OCD.

ABSTRACT

A program is presented to compute fallout fission product absorption in an expanding, cooling, uniform field of gas and fallout particles where the rate of fission product absorption is controlled by fission product surface concentrations as given by Henry's law constants and diffusion of these ission products into the fallout particles. The calculations are made a nuclide chain at a time, employing nuclear device and fission product parameters. Program output includes average concentration versus particle size for each absorbed fission product at a preselected lower temperature and the amount of that fission product in the gas phase, with an option for the calculation of fission product radial profiles versus particle size. As a demonstration of this model, an initial computer study of fallout formation phenomenology is presented.

CONTENTS

ABS	STRA	CT
SUN	MAR	Σ
1.	INTR	ODUCTION
2.	GENE	ERAL DESCRIPTION
3.	MET	HOD
4.	SUBP	PROGRAM DESCRIPTIONS
		Subroutine INITAL
		Main Routine NORMAN
		4.2.1. Setup and Calculations
		4.2.2. Main Loop
		4.2.3. Output
	4.3.	Function DIFFUS (DT, R)
	4.4.	Subroutine FINAL (AVCON, GAS, R, WT, NAME) 24
	4. 5.	
	4.6.	Function ERF (X)
		Subroutine BCALC (N. NT. I. KLAG, B. DD. DT) 28
API	PEND	ICES
	A.	Derivation of Equations Used
		Model Calculations
		Control Card Formats
		Sample Output and Subprogram Listings 59
REI	ERE	NCES

FIGURES

1. S	tep function approximation to time-temperature distribution 5
2. F	low chart, main loop NORMAN
A. 1.	Decay paths for n isotopes through N time steps
B. 1.	Calculated absorption in fallout of members of 137 chain in g fission product per g fallout as a function of temperature and particle size
B. 2.	Calculated absorption in fallout of members of 95 chain in g fission product per g fallout as a function of temperature and particle size
B. 3.	Concentration profiles at 900°K of 95 nuclide chain members as a function of particle size for standard detonation 48
B. 4.	Tellurium-137 concentration profiles at 900°K for standard particle size distribution and standard detonation 49
B. 5.	Iodine and xenon-137 concentration profiles at 900°K for standard particle size distribution and standard detonation 50
B. 6.	Cesium-137 concentration profiles at 900°K for standard particle size and standard detonation
B. 7.	Iodine-137 concentration profiles as a function of final temperature for the 0.0003 cm size particles and the standard detonation
	TABLES
1. V	ariables in common blocks
B. 1.	Standard particle size distribution
B. 2.	Concentration ratios at 900°K of particles of different size for various starting temperatures
B. 3.	Effect of size of temperature increment on fission product absorption
B. 4.	Effect of particle size distribution on fission product absorption
B. 5.	Variation of fission product absorption with size and height of detonation effects for a 100-KT fission yield detonation 47

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R. F. Korts and J. H. Norman General Atomic Report GA-7598

SUMMARY

The computer code described in this report represents a model of fallout fission product absorption, the fallout particles being distributed uniformly in a cooling, expanding field of gas. Rates of fission product absorption are governed by fission product surface concentrations as given by Henry's law constants and by diffusion into the fallout particles. It is assumed that the particles neither agglomerate nor convect fission products during the time stepped calculational portion of the program. The absorption of fission products is considered for one decay chain at a time using detonation yields, initial fission product yeilds, half-lives, Henry's law constants, concentration independent diffusivities, and a fallout particle size versus quantity distribution. Output consists of surface concentration versus temperature for each absorbed fission product, average concentration versus particle size at a final lower temperature, and the amount of that fission product remaining in the gas phase. Optionally, radial profiles of fission product concentration versus particle size at this temperature are calculated.

Fission product distribution for mass chains 95 and 137 in fallout using this code are described. One of the important findings derived from these computer studies is an understanding of how this model considers fractionation to vary with fallout particle size distribution.

I. INTRODUCTION

Since the phenomenological discovery of fractionation of fission product elements in recovered fallout, the importance of the chemical properties of the fission products and fallout particles to the description of fallout formation has been recognized. In his treatise on fallout, Miller (1) has attempted to synthesize fractionation effects by allowing fission products to equilibrate with fallout until they have cooled to 1400°C and, thereafter, to surface condense on the fallout. This scheme has been considered an oversimplification and phenomenologically unsatisfactory. General Atomic has taken the position that condensed state diffusion of fission products in glassy silicate fallout particles can be considered to control the rate of fission product absorption. This report is a description of a calculational program written in FORTRAN IV for the UNIVAC 1108 in which the surface of fallout is considered to be in equilibrium with the gas phase, but bulk fallout phase absorption of fission products is diffusion controlled. In addition, fallout formation information derived from some computer experiments employing this model is presented in Appendix B.

2. GENERAL DESCRIPTION

The problem consists of determining concentrations of radioactive isotopes in fallout particles. Absorption of fission products is considered a decay chain at a time where a chain consists of a number of isotopes generally of the same mass. The last isotope in the chain is assumed to be stable for the time periods being considered. The fallout particles are assumed to be non-agglomerating, internally non-convecting, refractory, spherical glassy silicate bodies of various sizes uniformly distributed throughout a uniform temperature-concentration cloud of particles, air, and fission products. Immediately after the detonation takes place, temperatures are very high and chemistry in the cloud is unimportant. However, at a time when the cloud has cooled sufficiently to consider chemistry, this program can be applied. From this point, as time progresses (and the temperature of the cloud further decreases) the isotopes are assumed to surface condense according to Henry's law and subsequently to be diffused into the particles. Concurrently, the isotopes decay from one to another at rates dependent upon their respective half-lives. At a sufficiently low temperature, the diffusion process is considered to terminate leaving the isotopes distributed among and within the particles in some manner. The purpose of the program is to describe that distribution in terms of surface co. centrations, average concentrations, and concentrations throughout the particles (profiles). To find these various concentrations, it suffices to calculate gas and condensed phase mass balances in which surface concentrations of each isotope are determined as a function of time.

Time-temperature stepping is employed in the program, and mass balances are determined at each time step. In order to determine the mass

balances, it is necessary to use input nuclear cloud parameters and chemical information. Time scaling functions for temperature and for cloud volume and a fallout mass representation are employed according to Miller⁽¹⁾ but can be changed to fit other circumstances. Fission product yields can be taken from Crocker, ⁽²⁾ and fission product decay rates can be taken from Crocker, Scheidt, and Connors. ⁽³⁾ Henry's law constants for this program have been presented by Norman. ⁽⁴⁾ A similar report on diffusion coefficients for this program is planned.

Surface concentrations are obtained directly from the mass balance calculations. Other output parameters are calculated from the surface concentration values and the degree of diffusion that has been calculated. The computation of concentration profiles of the various sized particles at a final temperature is optional. Another option allows calculation of late distributions at such a time that all remaining gaseous short-lived isotopes have decayed and condensed. In this option, it is assumed that the isotopes are deposited on the particles according to surface area.

Surface concentrations are calculated at a finite number of points assuming that an actual time-temperature distribution can be closely approximated by a step function as shown in Fig. 1. The exact nature of this step function depends on the starting, final, and incremental temperatures (see Section 4.1, item 4). Thus, all temperature (and time) dependent quantities are assumed to be step functions as well; in particular, volumes and diffusion constants and Henry's law constants for each isotope remain constant for each temperature interval. Fission product surface concentrations are also assumed constant during any time interval.

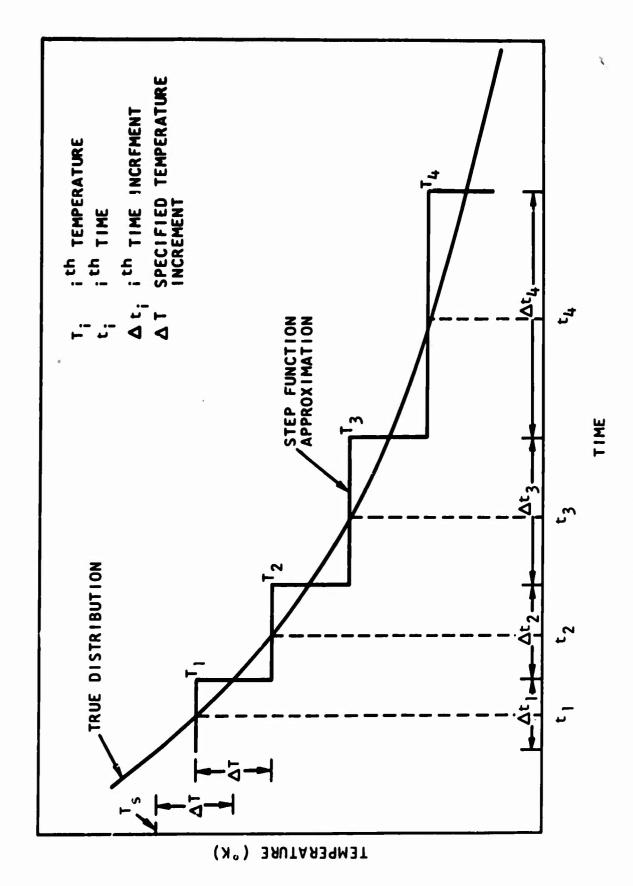


Fig. 1 -- Step function approximation to time-temperature distribution

3. METHOD

To determine the surface concentration of a given isotope at some temperature, the surface concentrations of that isotope and all its precursors at all preceding temperatures must be employed, since some fraction of each precursor will decay into the next isotope and previous surface concentrations will affect the total surface concentration at the specified temperature. We start by assuming that all of the fission products are uniformly distributed throughout the cloud at some sufficiently high temperature. To find surface concentrations at the first temperature, the following expression is used:

$$\frac{C_{i1}V_{1}M_{i}H_{i1}}{RT_{1}} + C_{i1}\sum_{s=1}^{p}D_{i11}^{(s)}w_{s} = Y_{i1}$$
 (1)

where C_{il} = surface concentration $(g/g)^*$ of isotope i at time t_i (sec)

V₁ = volume of nuclear cloud at time t₁ (liters)

M_i = molecular weight of isotope i (g/mole)

 H_{il} = Henry's law constant for isotope i at time t_{l} (atm/g/g)

R = ideal gas constant (liter atm/^oK mole)

T₁ = temperature at time t₁ (oK)

Y; = yield, or total amount of isotope i, available at time t, (g)

w = weight of soil (g) contained in particle size fraction s of
 radius r_s (cm)

 $D_{ill}^{(s)}$ = solution of the diffusion equation as if diffusion had taken place at a constant surface concentration C_{il} (see Section 4.3) for particle of sth radius, with the $\delta_{il}\Delta t_{l}$ term for isotope i assumed constant during Δt_{l} (δ_{il} is the pertinent diffusion coefficient (δ in cm²/sec)). The diffusion coefficients are assumed to be concentration independent.

^{*}Suggested self-consistent units.

The Y_{i1} value is computed by considering the decay of isotopes and the H_{i1} value is computed from characteristics of the isotope and temperature (see Section 4.1).

The first term on the left of Eq. (1) represents the amount of isotope i in the gas phase at time 1. The second term is the amount in the solid, or condensed, phase. For each s,

$$D_{i11}^{(s)} = \frac{\overline{C}_{i1}^{(s)}}{C_{i1}}$$

for the sth particle size, where $\overline{C}_{il}^{(s)}$ is the average concentration during time t_1 . Thus

$$C_{il} \sum_{s=1}^{p} D_{ill}^{(s)} w_{s} = C_{il} \sum_{s=1}^{p} \frac{\overline{C}_{il}^{(s)} w_{s}}{C_{il}}$$

$$=\sum_{s=1}^{p} \overline{C}_{il}^{(s)} w_{s}$$

It should be noted that in this program an approximation system is used in representing the decay of fission products. During the time intervals used, a fission product is assumed to decay to one and only one product. Thus, decay is not treated as a dynamic process but as a discrete process occurring only at the end of a time step period. A nearly correct fission-product distribution is obtained at the initiation of the program by stepping the time in many small increments to this initial time after detonation. The time increments then become governed by the temperature increments. This system will cause some deviation from an exact decay solution but, generally, will only alter the final answers in a negligible fashion. Of course, as the temperature increment approaches zero, the decay system will approach the theoretical limit.

To find surface concentrations at the second temperature, the expression is

$$\begin{split} &\frac{C_{i2}V_{2}M_{i}H_{i2}}{RT_{2}} + C_{i1} \exp(-k_{i}\Delta t_{1}) \sum_{s=1}^{p} D_{i,1,2}^{(s)} w_{s} \\ &+ C_{(i-1)1} \left[1 - \exp(-k_{i-1}\Delta t_{1})\right] \sum_{s=1}^{p} D_{(i-1),1,1;i,2,2}^{(s)} \\ &+ \left\{C_{i2} - C_{i1} \exp(-k_{i}\Delta t_{1}) - C_{(i-1)1} \left[1 - \exp(-k_{(i-1)}\Delta t_{1})\right]\right\} \sum_{s=1}^{p} D_{i,2,2}^{(s)} w_{s} = Y_{i2} \end{split}$$

The terms on the left side of Eq. (2) can be explained as follows:

1. The first term is due to the amount of isotope i still in the gas phase.

(2)

- 2. The second term is due to the amount of i surface condensed associated with time t_1 decaying for Δt_1 and being diffused for $\delta_{i1}\Delta t_1 + \delta_{i2}\Delta t_2$.
- 3. The third term is due to the amount of isotope i-1 (the immediate precursor at t_1 of i) surface condensed at time t_1 decaying into i and being diffused as i-1 during time Δt_1 and as i during time Δt_2 .
- 4. The fourth term represents the amount of i being diffused within the particles as the result of the surface concentration increment caused by the temperature-time-volume increment.

The D's are further defined in Appendix A.

It is apparent that as time is increased, the number of terms in the expression for C_{iN} becomes quite large. In fact, it can be shown (Appendix A) that if there are n-1 predecessors to an isotope, the number of terms required to find C_{iN} is

$$2^{N} - 1$$
 if $N \le n$

and

$$\sum_{k=0}^{n-1} \left(\sum_{j=n}^{N-1} {j \choose k} \right) + 2^{n} - 1 \quad \text{if } N \ge n$$
 (3)

plus additional terms for the gas phase contribution and yields.

The general expression for surface concentration can be written as

$$C_{iN} \frac{M_i V_N^H_{iN}}{RT_N} + A_{iN} = Y_{iN} , \qquad (4)$$

where A_{iN} consists of the number of terms described by Eq. (3). Each term of A_{iN} is of the form

$$C_{(i-r)(N-j)}$$
 (decay term) $\sum_{s=1}^{p} (D_{(1)} - D_{(2)}) w_s$, (5)

where $0 \le r \le i-n+1$

$$0 \le j \le N-1$$

The decay term depends on the position of the isotope i-r in the chain and on the time. $D_{(1)}$ and $D_{(2)}$ are diffusion terms which also depend on the position and time. (See Appendix A for details.)

Once surface concentrations have been found for all times, it remains to compute average concentrations and concentrations at various radii within the particles. Average concentrations can be obtained by mathematically combining concentration increments according to an additivity rule, as explained in Appendix A. Concentrations at any point within a particle can be found in an analagous way using relationships given by Crank. (5)

4. SUBPROGRAM DESCRIPTIONS

This section describes each of the FORTRAN IV subroutines used by this code. Most of the variables used are stored in one of the common blocks described in Table 1.

4. 1. SUBROUTINE INITAL

This routine reads the data cards for each calculation and prepares the input data for the main program. The flow is in approximately the order described below.

- Read input cards. The following information is read for each calculation with the internal variable name indicated in each case. (See Appendix C for card formats.)
 - a. Length of the chain, LC (<6)
 - b. Yield of device used for calculating fission product yields (in kilotons TNT), YKT
 - c. Yield of device used for calculating the time-temperature distribution and the cloud volume (in kilotons TNT), BKT
 - d. For each isotope in the chain:
 - (1) Chemical symbol, NAME
 - (2) Molecular weight (in g/mole), WMOL
 - (3) Initial yield (in atoms/10⁴ fissions), YI
 - (4) Half-life (in seconds), HL
 - (5) Constants used in calculating Henry's law constants and diffusion coefficients (see item 8), DC1, DC2, HC1, and HC2
 - e. Starting temperature, final temperature and incremental temperature (in 0 K), HTMP, ETMP, and TEMPV (also referred to as T_{g} , T_{f} , and ΔT)

TABLE I VARIABLES IN COMMON BLOCKS

Variable	Dimensions	Description				
BLANK Common: Contained in NORMAN, BCALC, INITIAL						
YLD	(6, 40)	Yield of each isotope at each time step				
NAME	(6)	Chemical symbol of each isotope				
WMOL	(6)	Molecular weight of each isotope				
DFC	(6, 40)	Diffusion coefficients of each isotope at each time step				
HEN	(6, 40)	Henry's law constants for each isotope at each time step				
AK	(6)	Exponential constant associated with each half-life				
R	(30)	Radius of each particle size				
WT	(30)	Weight of soil in each particle size				
VOL	(40)	Volume of the cloud at each time step				
T	(40)	Time increments				
TEMP	(40)	Temperatures at each time step				
ATM	(40)	Times at each step				
SII	NG: Contained	in NORMAN, FINAL, BCALC, INITIAL				
LC	(1)	Length of chain				
IR	(1)	Number of particle sizes				
NTOT	(1)	Number of time steps				
NCSI	(1)	Case identification				
NCS2	(1)	Case identification				
TWT	(1)	Total mass of soil				
ILG	(1)	Flag: if = 0, do not do FINAL calculation if ≠ 0, do FINAL calculation				
IDPTH	(1)	Number of radii used for profile calculations				
Q	(11)	Relative radii for profile calculations				
B1: Contained in NORMAN, BCALC, DELTA						
DCY	(6, 40, 40)	Array of decay constants for each isotope from each time step to each other				
WN	(6, 40)	Array of concentration ratios times weights for each isotope at each time step				
INLOOP: Contained in NORMAN, BCALC, DELTA						
L	(1)	Subscript of isotope presently under consideration				

- f. Number of particle sizes, IR (≤30)
- g. For each particle size, its radius (in cm) and the fraction of soil contained in particles of that size, R. PW
- h. The total mass of soil (in g/KT), TWGT
- 2. Check starting temperature. If the starting temperature is so high that all the fission products are uniformly distributed throughout each particle, that temperature is lowered by one increment and the procedure is repeated until a temperature is reached where this is not the case. Calculations are initiated after stepping back one temperature increment or to the starting temperature. To accomplish this, a D for the largest particle radius is found for each isotope at an initial time step. A D of I implies that an isotope is uniformly distributed throughout all particles. If D = I for all isotopes and particle sizes, the process is not yet diffusion controlled; thus, if

$$\sum_{i=1}^{LC} D_i = LC ,$$

all D's are unity. (Note that $D \le 1$.) If this is satisfied, the temperature is reduced and the procedure is repeated.

- 3. Calculate w_g , the weight of soil in each particle size: $w_g = TWT \cdot PW_g$, where $TWT = TWGT \cdot BKT$.
- 4. Calculate the time-temperature distribution. Using the formula (2)

$$T = 4.66 \times 10^3 (BKT)^{-0.010} exp(-0.546(BKT)^{-0.373}t) = f(t)$$
, (6)

times and incremental times are calculated to correspond to each specified temperature. The temperatures are calculated from

$$T_n = T_s - \left(\frac{2n-1}{2}\right) \Delta T ,$$

and the times such that

$$T_n = f(t_n)$$
.

 Δt_n is defined by

$$\Delta t_n = f^{-1} \left(T_n + \frac{\Delta T}{2} \right) - f^{-1} \left(T_n - \frac{\Delta T}{2} \right) \text{ (see Fig. 1)} . \tag{7}$$

5. Convert initial yields, Y_i, to grams.

$$Y_i(g) = YI \cdot WMOL \cdot YKT \cdot 2.3410261 \times 10^{-5}$$

 Find the quantities k_i. These quantities are used to calculate decay coefficients where

$$d_{inm} = \exp[-k_i(\Delta t_n + \Delta t_{n+1} + \dots + \Delta t_m)] \qquad m > n$$

$$= \exp(k_i \Delta t_n) \qquad m = n$$

= 1 (a computational aid) m < n

and

$$k_i = \frac{\ln(2)}{HL_i}.$$

7. Calculate yields at t_{g} where

$$t_{s} = f^{-1}(T_{s}) .$$

Yields at t are approximated by the following scheme:

Let $\Delta h = \min(HL)/10$ when $\min(HL)$ is the minimum half-life in the chain

Let $\tau_n = n (\Delta h)$, and let i-1 specify the precursor of isotope i. Then,

$$Y_{i}(\tau_{1}) = Y_{i}(0) \exp(-k_{i}\Delta h) + Y_{i-1}(0) [1 - \exp(-k_{i-1}\Delta h)]$$
 (8)

.

$$Y_{i}(\tau_{n}) = Y_{i}(\tau_{n-1}) \exp(-k_{i}\Delta h) + Y_{i-1}(\tau_{n-1})[1 - \exp(-k_{i-1}\Delta h)]$$
 (9)

Equation (9) is repeated for all nand i until n is such that $\tau_{n+1} > t_s$. For the last step, Δh is changed such that $\Delta h = t_s - \tau_n$ and Eq. (9) is repeated. If the original $\Delta h \ge t_s$, Δh is set equal to t_s and only Eq. (8) is used.

8. Calculate diffusion coefficients and Henry's law constants.

The expressions (4)

$$\log(H_{in}) = HCl_i - HCl_i / T_n$$
 (10)

and

$$\log(\delta_{in}) = DCl_i - DCl_i / T_n$$
 (11)

are used to find δ and H for all isotopes and temperatures.

9. Volumes are computed for all temperatures using the expression (3)

$$V_n = \left\{5.69 \times 10^2 (BKT)^{1/3} \exp[0.17(BKT)^{-0.373} t_n]\right\}^3 \frac{4\pi}{3}$$

10. Input and computed quantities are printed. Input is actually printed as it is read and calculated quantities as they are calculated. (See Appendix D for sample printout.)

4. 2. MAIN ROUTINE (NORMAN)

The main routine has three basic functions:

- 1. Setup and initial calculation of arrays needed (such as yields and decay coefficients) for the main loop.
- 2. Execution of the main loop which finds surface concentrations, average concentrations, and profiles.
- 3. Writing out of results.

4. 2. 1. Setup and Initial Calculations

The main routine starts by calling subroutine INITAL, which supplies it with necessary input data. Next it computes the array of decay coefficients, dinm, as defined in Section 4.1, item 6; the array of

yields Y , where Y is the yield of isotope i at time n; and the arrays UN and WN. The yields are found by using Eq. (8), Section 4.1, item 7, modified as

$$Y_{in} = Y_{i,n-1}d_{i,n-1,n-1} + Y_{i-1,n-1}(1 - d_{i-1,n-1,n-1})$$
.

The UN and WN arrays are defined by

$$UN_{ine} = E(\delta_{in}\Delta t_{n}, r_{e})$$

$$WN_{in} = \sum_{n=1}^{p} UN_{ins}^{w}$$

 δ_{in} = diffusion coefficient of isotope i at time t_n $\Delta t_n = n^{th}$ time interval

r = radius of s th particles

w = weight of soil in the sth particle size

 $E(\delta_{in}\Delta t_n, r_n) =$ solution of diffusion equation for sphere of radius r_n and

Finally, the surface concentrations are calculated for each isotope at time t, by the relation

$$C_{il}\left(\frac{WMOL_{i}H_{il}V_{l}}{RT_{1}}+WN_{il}\right)=Y_{il}$$

4. 2. 2. Main Loop

The following steps are taken in the main loop for each isotope and each time:

- 1. Find the terms due to decay of this isotope from previous times.
- 2. For each precursor of this isotope at each preceeding time, find all possible paths for decay, and for each such path find a diffusion term (see Section 4.8) and a decay term (see Section 4.7),

^{*}UN_{ins} = D_{inn} See Fig. 2

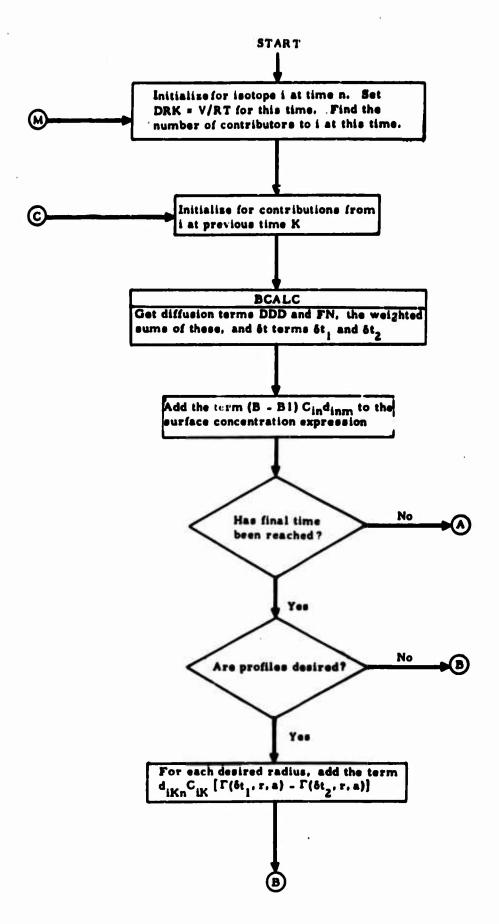


Fig. 2--Flow chart, main loop NORMAN (Sheet 1 of 5)

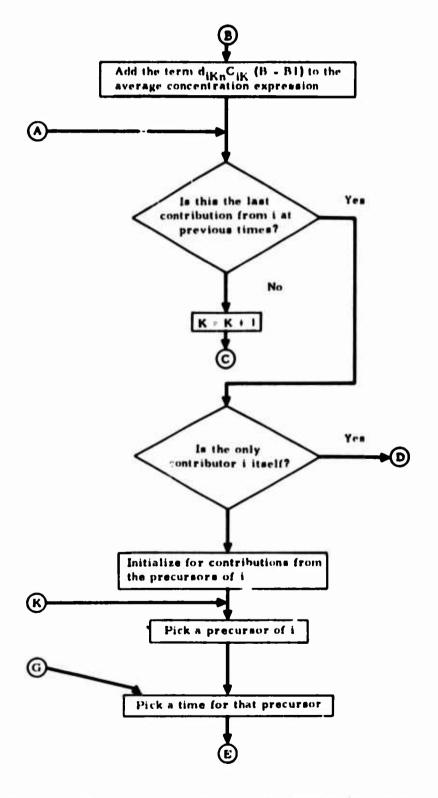


Fig. 2-- Flow chart, main loop NORMAN (Sheet 2 of 5)

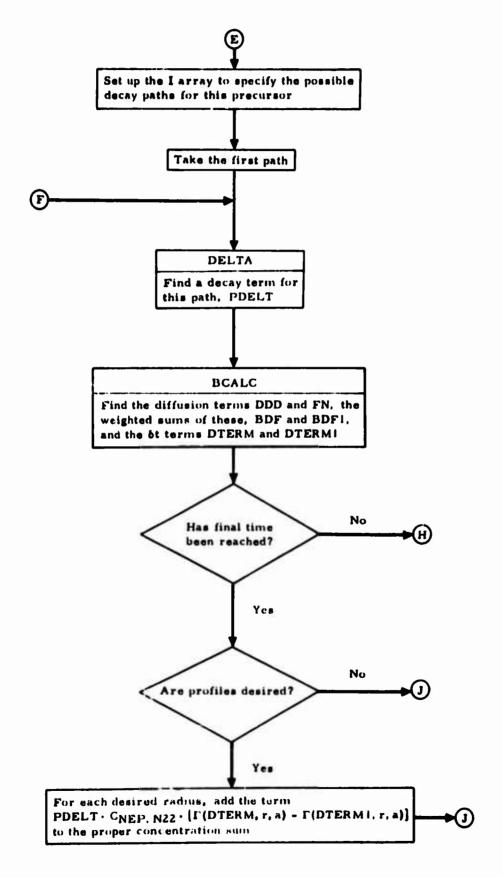


Fig. 2--Flow chart, main loop NORMAN (Sheet 3 of 5)

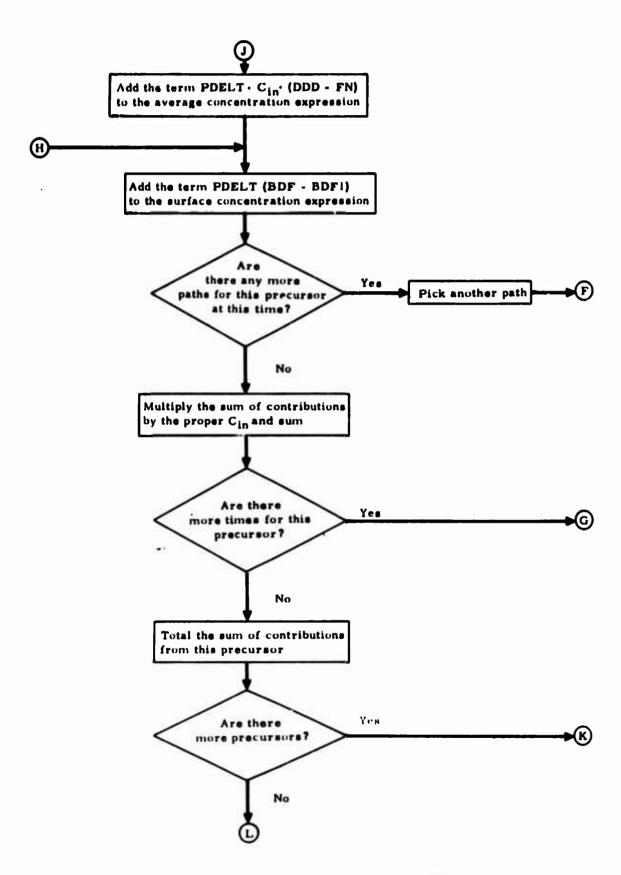


Fig. 2--Flow chart, main loop NORMAN (Sheet 4 of 5)

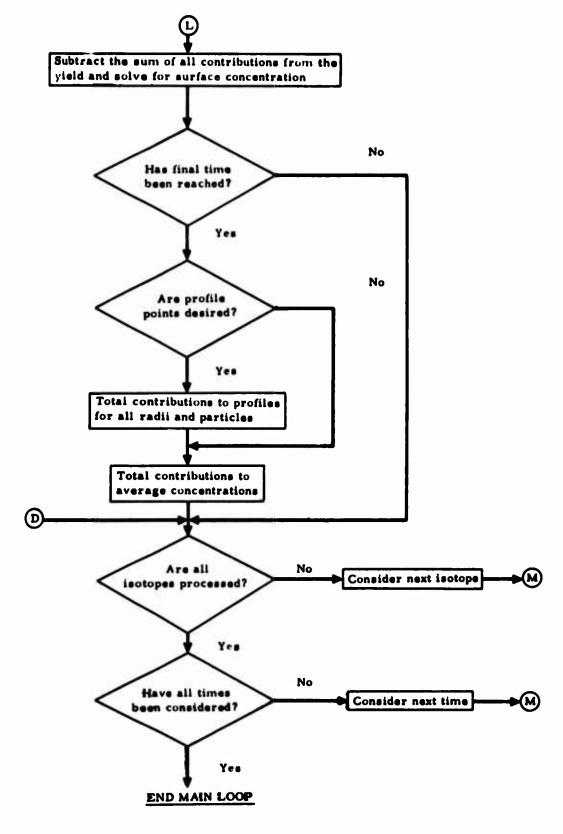


Fig. 2--Flow chart, main loop NORMAN (Sheet 5 of 5)

forming the product as in Eq. (5) and adding to the previous summed contributions. In the event that the final time step has been reached, add the appropriate terms to the average concentration and profile point sums.

 Let α represent the total of the sum of all contributions to the average concentration expression except that term due to absorption diffusion of this isotope from surface concentration C_{in}. Then the surface concentration is found from

$$C_{in}\left(\frac{WMOL_{i}H_{in}V_{n}}{RT_{n}}+WN_{in}\right)=Y_{in}-\alpha.$$

4. 2. 3. Output

Certain other calculations, as follows, are made while printing out results.

 Calculation of the amount of material remaining in the gas phase for each isotope at each time:

$$GAS_{in} = \frac{C_{in}H_{in}V_{n}WMOL_{i}}{RT_{n}}$$

2. Computation of the amount of material in the solid phase (diffused), per isotope, at the final time:

$$SOLID_i = \sum_{s=1}^p A_{is} w_s$$
,

where A is the average concentration of isotope i in particle size s.

3. Computation of the weighted mean average concentration for each isotope i:

$$\overline{A}_{i} = \frac{SOLID_{i}}{(total weight)}$$

The following quantities are printed:

a. T_n , t_n , and C_{in} for each i, n

b. T_n, t_n, and GAS_{in} for each i, n

c. T_N, r_s, and A_{is}, N being the final time and for each i, s

d. A, for each i

e. C_{is}(r) for each r, i, s, where C_{is}(r) is the concentration of isotope i in particle s at radius r

f. A summary of the amounts of material in gas and solid phases at the final temperature

4. 3. FUNCTION DIFFUS (DT, R)

This function computes a solution to the diffusion equation for a spherical body of radius r. The true solution is: (5)

$$D = 1 - \frac{6}{\pi^2} \sum_{n=1}^{\infty} \frac{1}{n^2} \exp(-n^2 \pi^2 \delta t / r^2)$$
 (12)

or

$$D = 6 \left(\frac{\delta t}{r^2}\right)^{1/2} \left(\pi^{-1/2} + 2 \sum_{n=1}^{\infty} ierfc \frac{nr}{\sqrt{\delta t}}\right) - 3 \left(\frac{\delta t}{r^2}\right) , \qquad (13)$$

where (6)

$$ierfc(x) = \int_{x}^{\infty} erfc(u)du$$

and

$$\operatorname{erfc}(x) = \frac{2}{\sqrt{\pi}} \int_{x}^{\infty} e^{-u^2} du$$
.

^{*}See Appendix D for sample output.

This printout and the calculation are optional (see Section 2).

(See Appendix A for details.)

Function DIFFUS uses the following approximations to these solutions:

Let
$$x = r/\sqrt{\delta t}$$

1. If x > 3, 14.

$$D = \frac{3}{x} \left(\frac{2}{\sqrt{\pi}} - \frac{1}{x} \right) .$$

which is a truncation of Eq. (13) with the infinite series neglected.

2. If 2. $0 < x \le 3.14$,

$$D = 1 - \frac{6}{\pi^2} \left[\exp(-\pi^2/x^2) + \frac{1}{4} \exp(-4\pi^2/x^2) \right]$$

3. If 0. 745 < $x \le 2$. 0,

$$D = 1 - \frac{6}{\pi^2} \left[\exp(-\pi^2/x^2) \right]$$

4. If $x \le 0.745$,

$$D = 1$$
.

where 2, 3, and 4 are truncations of Eq. (12).

The errors incurred by these approximations are bounded as follows:

$$|E_4| \le 2.5 \times 10^{-8}$$

 $|E_3| \le 1 \times 10^{-5}$
 $|E_2| \le 1 \times 10^{-5}$
 $|E_1| \le 5 \times 10^{-6}$

where E_i is the error using approximation i.

4. 4. SUBROUTINE FINAL (AVCON, GAS, R. WT, NAME)

Subroutine FINAL computes the concentration of the last isotope in the various particles assuming that

- 1. All other isotopes have completely decayed into the last one
- 2. None of the last isotope has decayed away

 All of the material in the gas phase when the main process was terminated has condensed onto the particles according to surface area

Let $a_{is} = average$ concentration of isotope i in particle of size s, $Z_{s} = r_{l}/r_{s}, \text{ where } r_{l} \text{ is the radius of the largest particles, and}$ $g_{i} = amount \text{ of isotope i in the gas phase when main process was}$ terminated.

Then,

$$f_s = \sum_{i=1}^{LC} a_{is}$$

$$G = \sum_{i=1}^{LC} g_i$$

 $V_{g} = Z_{g} w_{g}$ (proportional to the surface area)

$$v = \sum_{s=1}^{p} V_{s}$$

where w = weight of particles of size s.

Also.

$$F_{s} = \frac{Z_{s}G}{v} + f_{s}$$

and

$$W_s = F_s W_s$$
 ,

where $F_s = \text{total average concentration (includes gas condensation) in particle of size s$

W = total weight of isotope in particle of size s

 f_{s} = average concentration in particle of size s due to the decay of solids.

The average concentration of the last isotope of the chain, f_s , is considered to be equal to the sum of the average concentrations of all the members of the chain when the main process was terminated. Subroutine FINAL also lists out f_s , and f_s for each particle size.

4.5. FUNCTION PROFIL (DT, r,Q)

Function PROFIL computes the concentration at a distance a = rQ from the center of a spherical particle of radius r. The value of the concentration is found from the expression: (5)

$$\frac{C}{C_o} = 1 + \frac{2r}{\pi a} \sum_{n=1}^{\infty} \frac{(-1)^n}{n} \sin \frac{n\pi a}{r} e^{-n^2 \pi^2 \delta t/r^2}$$
 (14)

or

$$\frac{C}{C_o} = \frac{r}{a} \sum_{n=0}^{\infty} \left\{ \text{erfc} \left[\frac{(2n+1)r - a}{2\sqrt{\delta t}} \right] - \text{erfc} \left[\frac{(2n+1)r + a}{2\sqrt{\delta t}} \right] \right\}, \quad (15)$$

where r = radius of particle

a = distance from center of particle, $a \ge 0$. lr

C = surface concentration (assumed constant)

Function PROFIL computes the value of the ratio C/C using a truncation of Eq. (14) or Eq. (15). In all cases, the series is terminated when either

$$\frac{|S_n - S_{n-1}|}{|S_n|} < \epsilon^*$$

or when n > 100, where S_k is Eq. (14) or Eq. (15) truncated to k terms.

The series used is chosen using the following criteria:

If $\delta \Delta t < b_j$, Eq. (15) is used where j is determined as the smallest integer such that

$$r \ge 10^{-j}$$
,

where $b_i = 10^{1-j-[(j-1)/2]}$

and where [x] = greatest integer I such that $I \ge x$.

Otherwise Eq. (14) is used.

 $[\]epsilon = 0.01$ at present.

This scheme was determined by examining the values of Eqs. (14) and (15) for various numbers of terms and picking the equation that required the least number of terms.

4.6. FUNCTION ERF(X)

Function ERF computes an approximate value of the error function using a rational approximation. The error function is defined as:

$$erf(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-t^2} dt .$$

ERF uses the approximation (5)

$$erf(x) = 1 - (a_1t + a_2t^2 + a_3t^3)e^{-x^2}$$

where

$$t = \frac{1}{1 + px}$$

and

$$p = 0.47047$$
 $a_1 = 0.3480242$
 $a_2 = -0.0958798$
 $a_3 = 0.7478556$

The error E in this approximation is bounded by

$$|\mathbf{E}| \leq 2.5 \times 10^{-5} .$$

4. 7. SUBROUTINE DELTA (N, NT, I, PR)

Subroutine DELTA computes the decay coefficient, PR, associated with decay of isotope i to the presently considered isotope at time NT.

The array I specifies the path being taken, with I(1) specifying the time at which i is considered.

A product of terms of the form

$$d_{inm}(1 - d_{i, m+1, m+1})$$

is found. The d array is initially calculated in the main routine in such a way that

$$d_{inm} = 1$$
 if $n > m$.

The general product then has the form

$$PR = d_{inm}(1 - d_{i,m+1,m+1})d_{k,m+2,\ell}(1 - d_{k,\ell+1,\ell+1}) \cdots d_{n,r,NT-1},$$

where i is the subscript of the isotope presently under consideration. The values of n, ℓ , and r depend on the path.

4. 8. SUBROUTINE BCALC (N. NT, I, KLAG, B. DD, DT)

This routine computes the correct δt term associated with the path specified by I and calls DIFFUS to solve for a diffusion term.

The δt term described by Eq. (A-10) is computed from the relation

$$\delta t = \sum_{p=n}^{m} \delta_{ip} \Delta t_{p} + \sum_{p=k}^{\ell} \delta_{(i+1)p} \Delta t_{p} + \ldots + \sum_{p=v}^{q} \delta_{(i+r)p} \Delta t_{p},$$

where n, m, k, ℓ , ..., v, and q are specified by the I array. The corresponding diffusion terms,

are then found by calling DIFFUS for each s. Depending upon KLAG, either the first D or the second D in Eq. (A-9) is found and stored in the DD array. δt is stored in DT and the expression

$$\sum_{s=1}^{p} (D_{i, n, m; i+1, k, \ell; ...; i+r, v, q}^{(s)})^{w}$$

is stored in B.

^{*}See Section 4.1, item 6 for definition of dinm.

APPENDIX A

DERIVATION OF EQUATIONS USED

A. 1. THE GENERAL TERM IN THE SURFACE CONCENTRATION EQUATION

An analysis of the case of a single stable isotope with no precursors, which eliminates the decay question, makes it easier to understand the general case. If Eqs. (1) and (2) in Section 3 are rewritten leaving out the isotope identification and letting $w_1 = 1$ and p = 1, the following equations are obtained:

$$C_1G_1 + C_1D_{11} = Y$$
 (A-1)

and

$$C_2G_2 + C_1D_{12} + (C_2 - C_1)D_{22} = Y$$
, (A-2)

where

$$G_{N} = \frac{V_{N}^{MH}N}{RT_{N}}.$$

Equation (A-2) can be rewritten as follows:

$$C_2G_2 + C_2D_{22} + C_1 (D_{12} - D_{22}) = Y$$
.

The expression at t₃ is

$$C_3G_3 + C_1D_{13} + (C_2 - C_1)D_{23} + [C_3 - (C_2 - C_1) - C_1]D_{33} = Y$$

or

$$C_3G_3 + C_1D_{13} + (C_2 - C_1)D_{23} + (C_3 - C_2)D_{33} = Y$$
.

The expression at t_A is

$$C_4G_4 + C_1D_{14} + (C_2 - C_1)D_{24} + (C_3 - C_2)D_{34} + (C_4 - C_3)D_{44} = Y$$
.

The expression at step N is then

$$C_N^G_N + C_1^D_{1N} + \sum_{k=2}^{N} (C_k - C_{k-1}^D)_{kN} = Y$$
 (A-3)

Equation (A-3) can be rewritten as follows:

$$C_N G_N + C_1 D_{1N} + \sum_{k=2}^{N} C_k D_{kN} - \sum_{k=1}^{N-1} C_k D_{k+1, N} = Y$$

or

$$C_N G_N + \sum_{k=1}^{N} C_k D_{kN} - \sum_{k=1}^{N-1} C_k D_{k+1, N} = Y$$

or, finally, as

$$C_N(G_N + D_{NN}) + \sum_{k=1}^{N-1} C_k(D_{kN} - D_{k+1,N}) = Y$$
 (A-4)

Next, consider the general expression with decay from precursors Figure A·l represents n isotopes at N time steps; the connecting lines represent the ways in which decay can occur. The numbers occurring with each state represent the number of ways in which that isotope at that time can decay to isotope l at t_N . It is apparent that if $N \le n$, the diagram is triangular and is a Pascal's triangle, the numbers representing binomial coefficients. Thus for $N \le n$, the total number of terms in the surface concentration expression is

$$2^{0} + 2^{1} + 2^{2} + \cdots + 2^{N-1} = 2^{N} - 1$$
.

For N > n, the diagram is a truncated Pascal's triangle. In this case, there are $2^n - 1$ terms from the purely triangular portion, plus

$$\sum_{k=0}^{n-1} \left[\sum_{j=n}^{N-1} \binom{j}{k} \right]$$

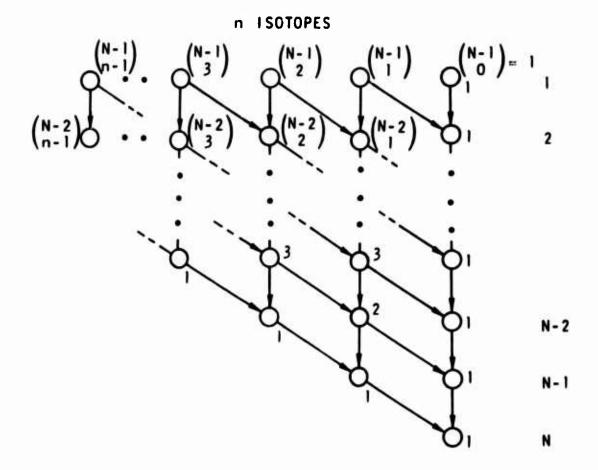


Fig. A. 1--Decay paths for n isotopes through N time steps

from the rectangular portion, giving a total of

$$M = \sum_{k=0}^{n-1} \left[\sum_{j=n}^{N-1} {j \choose k} \right] + 2^{n} - 1 . \qquad (A-5)$$

In Eq. (A-4), each term in the sum consists of a surface concentration times a diffusion term. In each case the diffusion term is the difference of two diffusion terms, one representing diffusion from t_k to t_N and the other from $t_{(k+1)}$ to t_N . Each of the M terms in the general expression is of a similar form, modified by a decay term.

As an example, consider the terms associated with the position n = 2in Fig. A.1 at t_A . There are three possible paths for decay to isotope 1. *The decay coefficients associated with these paths are

$$\begin{pmatrix} (1 - d_{211})d_{123} \\ d_{211}(1 - d_{222})d_{133} \\ d_{212}(1 - d_{233}) \end{pmatrix}$$
(A-6)

where dinmrepresents the decay coefficient of isotope i from time step Δt_n through Δt_m . The corresponding isotope indexed diffusion terms are:

$$\begin{array}{c}
D_{2,1,1;1,2,4} - D_{1,2,4} \\
D_{2,1,2;1,3,4} - D_{2,2,2;1,3,4} \\
D_{2,1,3;1,4,4} - D_{2,2,3;1,4,4}
\end{array}$$
(A-7)

where $D_{i,n,m;i+1,k,\ell;...;i+r,v,q}$ represents diffusion as isotope i from t_n to t_m , as isotope i+l from t_k to t_ℓ , etc., and as isotope i+r from time t_v to t_q (see Eq. (A-10)).

^{*}See Section 4.1, item 6.

Note that decay is carried only through Step 3, while diffusion is carried through Step 4 (see Section 3).

The total contribution to isotope 1 at t₄ from isotope 2 at t₁ can be written as

$$C_{21} [(1 - d_{211})d_{123}(D_{2, 1, 1; 1, 2, 4} - D_{1, 2, 4}) + d_{211}(1 - d_{222})d_{133}(D_{2, 1, 2; 1, 3, 4} - D_{2, 2, 2; 1, 3, 4}) + d_{212}(1 - d_{233})(D_{2, 1, 3; 1, 4, 4} - D_{2, 2, 3; 1, 4, 4})]$$
(A-8)

where again a single particle of unit weight is assumed.

From the above discussion it can be seen that the general term is of the form

 $C_{in}d_{inm}(1-d_{i,m+1,m+1})d_{i+1,m+2,k}(1-d_{i+1,k+1,k+1})\cdots$

···d i+r,
$$\ell$$
, N-1
$$\sum_{s=1}^{p} (D_{i,n,m+1;i+1,m+2,k+1;\cdots;i+r,\ell,N}^{(s)})$$

$$-D_{i,n+l,m+l;i+l,m+2, k+l; \cdots; i+r, \ell, N}^{(s)} w_{s}$$
 (A-9)

where

ľ

- We are considering the contribution to isotope i + r at time N
 via some path, from isotope i at time n.
- 2. $m \ge n + 1$ $k \ge m + 1$ \vdots $\ell \ge i + r$
- 3. There are p particle sizes with weights w_1, w_2, \ldots, w_p .
- 4. If n > m, $d_{inm} = 1$ as a mathematical convenience.

The quantity $D_{i, n, m; i+1, k, \ell; \cdots; i+r, v, q}^{(s)}$ represents a solution of the diffusion equation for a particle of radius r_s and δt term given by

$$\delta t = \sum_{n=q}^{n} \delta_{in} \Delta t_{n} + \sum_{n=k}^{\ell} \delta_{i+l,n} \Delta t_{n} + \cdots + \sum_{n=v}^{q} \delta_{l+r,n} \Delta t_{n} , \quad (A-10)$$

where δ_{in} is defined by Eq. (11), Section 4.1, item 8, and Δt_n is defined by Eq. 7, Section 4.1, item 4.

A. 2. AVERAGE CONCENTRATIONS AND PROFILES

As in the case of surface concentrations, the case considered is that of a single stable isotope with no precursors. The basic relationship is

$$\frac{\overline{C}}{C} = D , \qquad (A-11)$$

where \overline{C} is an average concentration increment. C is a surface concentration increment, and D is a diffusion term as in Section A. 1. At the first time step, it is clear that

$$\overline{C}_1 = C_1 D_{11}$$
.

At the second time step,

$$\overline{C}_2 = C_1 D_{12} + (C_2 - C_1) D_{22}$$

where the first term on the right represents the concentration at time step 1 further diffused and the second term represents material condensed and diffused during time step 2. Similarly, at time step 3,

$$\overline{C}_3 = C_1D_{13} + (C_2 - C_1)D_{23} + (C_3 - C_2)D_{33}$$
,

where the three terms represent, respectively, condensation during time steps 1, 2, and 3, and diffusion from those times to time 3. At time step N, it follows that

$$\overline{C}_{N} = C_{1}D_{1N} + \sum_{k=2}^{N} (C_{k} - C_{k-1})D_{k, N}$$
, (A-12)

where each term in the sum represents a concentration diffused over the specified time interval.

To obtain an expression similar to Eq. (A-4), Eq. (A-12) is rewritten

$$\overline{C}_{N} = C_{1}D_{1N} + (C_{N} - C_{N-1})D_{N, N} + \sum_{k=2}^{N-1} C_{k}D_{k, N} - \sum_{k=2}^{N-1} C_{k-1}D_{k, N},$$

or

$$\overline{C}_{N} = C_{N}D_{N,N} + \sum_{k=1}^{N-1} C_{k}D_{k,N} - \sum_{k=1}^{N-1} C_{k}D_{k+1,N}$$
,

or, finally, as

$$\overline{C}_{N} = C_{N}D_{N,N} + \sum_{k=1}^{N-1} C_{k}(D_{k,N} - D_{k+1,N}) . \qquad (A-13)$$

It can be seen that the general expression including decay from precursors consists of terms similar to Eq. (A-9). The number of such terms is the same as that described in Section A·1, since there is a contribution from each isotope at each preceeding time step for each possible decay path. Here the expression for average concentration in the shaper time size consists of terms of the form:

$$C_{in}^{d}_{inm}(1-d_{i,m+l,m+l})^{d}_{i+l,m+2,k}(1-d_{i+l,k+l,k+l})\cdots$$

··· d i+r,
$$\ell$$
, N-1 (D(s) i,n,m+l;i+l, n+2, k+l; ···; i+r, ℓ , N

-
$$D_{i,n+l,m+l;i+l,m+2,k+l;\cdots;i+r,\ell,N}^{(s)}$$
 (A-14)

summed over all contributions.

To find profiles, it is desired to find the concentration at specific points within the particle. The basic relation here is (see Ref. 6)

$$\frac{C(a)}{C} = \Gamma(a, r, \delta t) , \qquad (A-15)$$

where $\Gamma(a, r, \delta t) = 1 + \frac{2r}{\pi a} \sum_{n=1}^{\infty} \frac{(-1)^n}{n} \sin \frac{n\pi a}{r} \exp(-n^2 \pi^2 \delta t/r^2)$.

As with average concentrations, Eq. (A-15) gives a concentration increment due to diffusion during the given time and with zero initial concentration. For a stable isotope with no precursors we can, in an analogous manner, obtain the expression

$$C_{iN}(a) = C_{iN} \Gamma(a, r, \delta_{iN} \Delta t_{N}) + \sum_{k=1}^{N-1} C_{ik} \left[\Gamma\left(a, r, \sum_{n=k}^{N} \delta_{in} \Delta t_{n}\right) - \Gamma\left(a, r, \sum_{n=k+1}^{N} \delta_{in} \Delta t_{n}\right) \right]$$
(A-16)

and in the general case obtain a sum of terms similar to Eq. (A.14) with each D replaced by the appropriate Γ .

APPENDIX B MODEL CALCULATIONS

Some of the initial results of calculations obtained using this program are considered in this appendix. In these calculations examples of a refractory and a volatile chain, 95 and 137, respectively, were selected for study. The results should demonstrate the difference in behavior for extreme cases of refractory and volatile nuclide chains with varying conditions and thus provide limits of fission product behavior during fallout formation. Data sources for these calculations have been listed earlier. In addition to these sources, the standard particle distribution shown in Table B. I was developed from one used by Miller, experimental and estimated diffusion coefficients were employed as measured or estimated in reported studies, and total fissions were scaled as a function of weapon fission yield according to Miller.

Several program variables had to be investigated in order to demonstrate the degree of dependence of the calculational system on these variables. For instance, the variation of results with the size of the temperature increment must be vanishingly small for the selected temperature increment as smaller increments are considered. Also, the variation of results with choice of initial and final calculational temperatures should be defined. Finally, the fact that the particle size distribution employed is idealized—that is, some selected radii represent all of the mass of the particles—should be considered. These points have been investigated and seem to be satisfactorily understood as will be demonstrated as the calculations are presented.

The calculations for the 95 and 137 chains (neglecting I 137 \rightarrow Xe 136 and I 138 \rightarrow Xe 137) were performed using 2700°K as an initial temperature. 200°K increments, terminating at 900°K, and employing the particle distribution described in Table B. I. The size of the detonation was taken as 100 KT pure fission. The amount of soil associated with the cloud was 8.5 \times 10 9 g.

^{*}Miller. C. F., private communication, November 1965.

Table B. 1
STANDARD PARTICLE SIZE DISTRIBUTION

	Particle Radius	Weight Percent of
Designation	(cm)	Particle Fraction
T	0. 00030	0. 38
S	0. 0030	28. 92
M	0. 0154	60. 2
L	0. 13	10.5

The calculated descriptions of behavior of the 137 and 95 chains are presented in Figs. B. I and B. 2, respectively. In these figures, the average concentration of a particular nuclide in the chain is presented as a function of the last temperature for which the adsorption-diffusion process was calculated. Also included are the gas phase contents of the nuclides at the same temperatures (designated by open circles). In the phenomena described in Fig. B. I, the curve for the tellurium gas phase content is decreasing as the result of two processes:

- 1. Te¹³⁷ decay.
- 2. Absorption of Te¹³⁷ in the fallout particles.

Corresponding processes are involved to different degrees in the description of the gas phase content of any nuclide. In the case of Te¹³⁷ at the lowest temperature, so little of the Te¹³⁷ has been absorbed that the gas phase content curve closely represents the tellurium decay. The gas phase content curves for iodine and xenon also essentially represent radioactive decay phenomena alone, since both of these elements are very volatile. On the other hand, the gas phase content of cesium is dropping sharply with temperature, not because Cs¹³⁷ is decaying but, in this case, because 94% of the cesium in the system has been absorbed by the fallout particles by the 1900°K point and more is absorbed at lower temperatures. Thus, the cesium curve describes mainly the absorption process.

The average concentrations of Te¹³⁷ as they vary with temperature and particle size were also investigated. The large particles, L, as defined in Table B. 1, show a slow but continuous increase in average tellurium concentration as the temperature decreases (see Fig. B. 1).

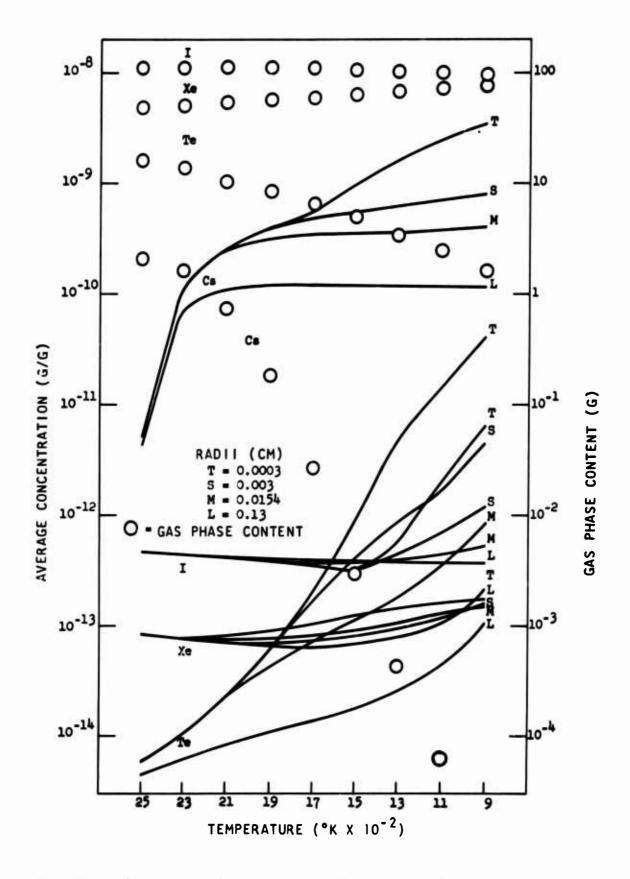


Fig. B. 1--Calculated absorption in fallout of members of 137 chain in g fission product per g fallout as a function of temperature and particle size

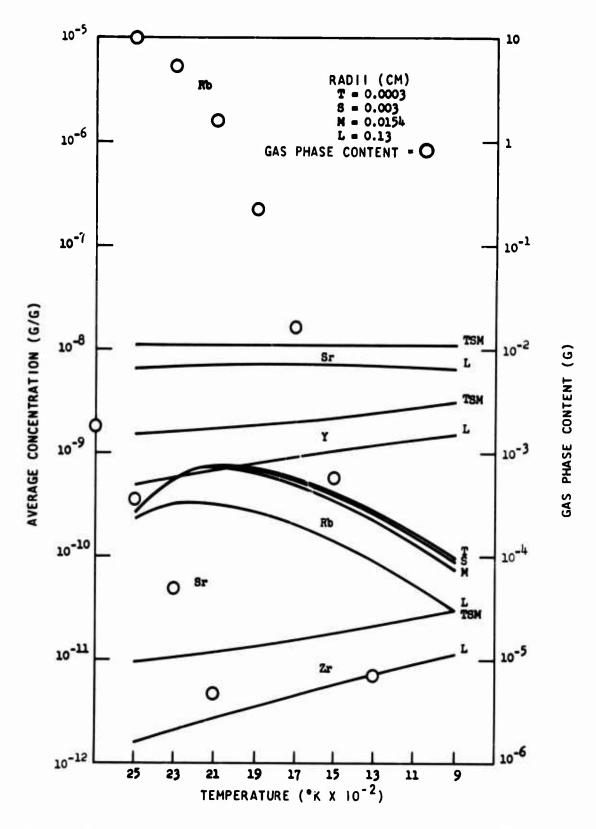


Fig. B.2--Calculated absorption is fallout of members of 95 chain in g fission product per g fallout as a function of temperature and particle size

Also, at the initial temperature the L particles exhibit a slightly lower concentration than the medium-sized particles, M. This point is important because it suggests that at the highest temperature of the calculation, the L particles did not equilibrate with the gas phase; that is, diffusion was not complete. In contrast, the tiny particles, T, the small particles, S, and the M particles exhibit essentially the same average concentration at the initial temperature. Further consideration of Fig. B. I reveals that the L particles experience a somewhat smaller increase in tellurium concentration than the M particles as the temperature decreases. Also, the S and M curves exhibit a departure from the more sharply rising T curve at two different lower temperatures. Departure points from the T curve occur as the central portions of the L. S. a 1 M particles effectively lose contact with the surface and diffusive transport begins to control the rate of absorption. In addition, at 1300°K absorption into the T particles is diffusion-controlled (at this temperature the surface concentration is calculated to be twice the average concentration).

The behavior of other elements in the 137 chain can also be observed in Fig. B.l. Neither iodine nor xenon is strongly absorbed at even the lowest temperature of this study. The higher temperature portions of the iodine and xenon curves in Fig. B.l reflect the low absorption and a low temperature coefficient of absorption for these isotopes. At lower temperatures, however, the iodine curves reflect the behavior of the parent tellurium curves. That is, at these temperatures the main iodine concentration in the particles comes from absorbed tellurium decay to iodine, whose loss from the particle is diffusion-controlled. The low temperature increase in xenon concentration in the particles is due to iodine decay, and xenon loss is also diffusion-controlled. This latter process has just become apparent as a result of the lowest temperature calculations shown in Fig. B.l. Similar phenomena can be observed in the cesium absorption curves. Fission cesium is mainly absorbed at 2100°K, where the absorption curves in Fig. B.l stop their steep ascent. This is the same

Note that the concentration of cesium in the L particles does not appreciably change at lower temperatures. The remaining gaseous cesium, and that formed in the gas phase decay of xenon, is picked up mainly by the smaller particles because of their large surface to volume ratio. The cesium concentration of the M particles does not change much after 1900°K, while that of the S particles begins to level off at 1700°K. The cesium concentration of the T particles increases sharply, reflecting the continued absorption of cesium produced from xenon (g).

The behavior of the 137 nuclide chain can thus best be explained in terms of diffusion-controlled sorption of tellurium and cesium and decay of the nuclides. Note that the sorption processes depicted in Fig. B. 1 account for only a minor portion of the total mass of fission products of this nuclide chain. The rest of the material, at times subsequent to those of this calculation, will go through the process of decaying to cesium and depositing on the outer portions of the fallout particles or on any condensed phase with which it comes in contact. In this study, however, it is assumed that the material remaining in the gas phase at 900°K is surface-deposited on the fallout particles themselves, according to their surface area. In reality, xenon, with a relatively long half life (234 sec), may well separate from portions of the fallout particles before decaying and depositing.

Figure B.2 shows similar data for the 95 chain. The behavior depicted in this figure, except for that of rubidium, represents essentially complete absorption at the highest temperature of the study, 2700° K, and rubidium absorption is essentially complete by 1900° K. The absorption differences between the smaller particles and the L particles result from the nuclides in this chain being incompletely diffused in the larger particles at 2700° K. The splitting of the M, S, and T curves of rubidium results from diffusion control becoming important at lower temperatures for smaller particles, as was discussed previously for tellurium in the 137 chain.

A further description of the effect of initial temperature is appropriate here. Table B.2 contains the 900°K average concentration ratios of L and M for zirconium, yttrium, and strontium and the L/M, M/S, and S/T average concentration ratios for rubidium. For a highly absorbed species, the ratio L/M differs from unity but obviously will approach unity as starting temperatures are increased. However, for rubidium, which is not so highly absorbed, the concentration ratios in various sized particles are independent of the starting temperature; and the smaller the particles, the more nearly this value approaches unity. The calculational values for less highly absorbed species thus appear to be independent of starting temperature.

Table B. 2

CONCENTRATION RATIOS AT 900°K OF PARTICLES OF DIFFERENT SIZE FOR VARIOUS STARTING TEMPERATURES

Starting Temperature		Ave	rage Conc	entration R	atio	
(°K)	Zr (L/M)	Y (L/M)	Sr (L/M)	Rb (L/M)	Rb (M/S)	Rb (S/Т)
2700	0. 378	0. 497	0.611	0. 398	0.844	0. 950
3100	0.600	0. 754	0.859	0. 398	0.844	0. 950
3900	0. 862	0. 924	ი. 883	0. 398	0.844	0. 950

The phenomenon of lack of equilibration of nonfractionated nuclides in the larger particles at high temperatures which is seen when the present model is used would seem to indicate some fractionation with size. An investigation of fallout should reveal whether this is indeed the case.

The orientation calculation described above has demonstrated many of the features of the diffusion-controlled model. These features seem to find some support in studies of fallout and certainly should be studied more thoroughly. However, before further questions concerning calculations of fallout phenomena are investigated, some of the strictly mathematical variations within the model calculations should be considered.

An investigation of different-sized temperature stepping intervals was made. In general, in the calculations reported here, 200°K temperature intervals were used. However, these test calculations, which were initiated at 2700°K and were carried through 900°K, were made for the two nuclide chains in both 200° and 100°K increments. Comparing results was more complicated than it first appeared because the final diffusion time corresponded to 800° K in the 200° interval calculations and to 850°K in the 100° interval calculations. In order to avoid this difficulty, both systems were allowed to decay and to absorb the resulting gaseous fission products (mainly Cs 137) according to surface area, thereby nearly eliminating the final calculational temperature differences. The comparison of results is presented in Table B.3. The only difference that seems significant occurred in the L particle zirconium concentrations, and it can apparently be explained by the difference in initial starting temperatures. The particles had essentially twice the time to diffuse at 2700°K in the case of the 200° increments as in the case of the 100° increments. While a better example might be constructed, it would appear from this study that the size of the temperature increment used in the model may not be very important if chosen as 200°K or less.

Table B. 3

EFFECT OF SIZE OF TEMPERATURE INCREMENT ON FISSION PRODUCT ABSORPTION

		Absor	ption(g)
Nuclide	Particle Size	100°C Temp. Interval	200°C (Std) Temp. Interval
Zr	T	0. 456	0. 456
	s	34.67	34.63
	M	71.52	71. 17
	L	6.86	7. 26
Cs	T	14.62	14.63
	S	112.40	112.44
	М	46. 92	46.88
	L	1.045	1.042

In the investigation of the effect of the particle size distribution, three different distributions were studied under the same conditions. The distributions employed and the weights of soil associated with each distribution in this case are presented in Table B.4. The amounts of cesium and zirconium fission products (in grams) after decay of precursor fission products are also presented for each fallout particle size as in the previous case. Finally, a ratio of the cesium and zirconium average concentrations associated with each particle size is presented. This ratio is proportional to a fractionation factor ratio comparing these elements. The proportionality constant is 0.649, a nuclide chain yield ratio. Further consideration of Table B. 4 reveals that fractionation factors for the very smallest particles far exceed unity, while for the larger particles these factors are less than unity. Fractionation factors in this model are thus quantities that depend on particle size (and also probably on the degree of detachment of the fallout particles from the nuclear cloud before conversion of volatile fission products to condensable ones). Therefore, if this model is correct, observed fractionation factors should be dependent on particle size and thus not accurately definable without ample description of this parameter. Indeed, these factors should be sensitive not only to the particle size in question but also to the particle size distribution. Specifically, one observes using distributions (2) and (3) that fractionation factors for the same size particles can differ considerably. Particle distribution (3) contains a large amount of very small particles. Distributions (1) and (2) are similar, since they both are based on Miller's description. Note that the distributions (1) and (2) produce similar fractionation factures for similarly sized particles. It seems apparent that in testing this model for reproduction of fractionation factors, one must specify the particle size distribution, a factor often not well documented.

A study was made of the effect of the size of the detonation and the amount of soil in the cloud (that is, the proximity of the detonation to the

Miller, C. F., private communication, November 1965.

Table B. 4

EFFECT OF PARTICLE SIZE DISTRIBUTION ON FISSION PRODUCT ABSORPTION

Radius Wei		Fallout (g)				Fission Product (g)	duct (g)				Ratio	
	Weight (1)	Weight (2)	Weish (2) Weish (3)	((,)	(2)	(,,)	7= (11)	7- (9)	7. (3)	\$ (1) = Z / 2 /	- () - () - () - ()	東(に) エム/ *) 東(に) エム/ *) (1) エム/ *)
+	†			;			11, 22	I	(6)	(1) 27.82	CB: 25 (5)	(6) 27,87
0 01			8. 5×10*			104.6			1.20			87.0
10-5		6.5	4.25×10 ⁵		0.078	52.7		1. 17×10-4	6.01		667.0	- eo
† .01		5. 95×10³	8. 5×10 ⁵		5. 50	0.11	Miller in P	0. 0816	12. 02		67.0	0. 92
3×10-4 3.23	3.23×104			14.6			0.444		-	33.0		
10-3		1. 19×10 ⁶	2. 975×10 ⁶		110.7	5. 13		. 16. 32	42.06	1	80	0. 12
3×10 ⁻³ 2.45	2.458×1C ⁶			112.4			33.73			3.3		
7-01	-	5.95×106	2.89×10 ⁶		57.3	1. 40		80.9	40.61		0.71	0.034
1.54×10 ⁻² 5.117×10 ⁶	901×21			46.9			69.28			0.68		
1.0		1.353×10 ⁶ 8.5×10 ⁵	8.5×10 ⁵		1.45	0. 132		16.2	10. 54		0.090	0.0125
25. 13 8. 72	8. 725×10 ⁵			3			10.06			0. 10		
1.0		8.415×10 ² 4.25×10 ⁵	4.25×10 ⁵		9.6×10-5	9.8×10-3		2.05×10 ⁻³	1. 07		0.046	0. 0092

²Can be converted to fractionation factors by multiplying by 0.649.

land surface). The calculated final amounts of cesium and zirconium for the various particle sizes as a function of these two quantities are shown in Table B. 5. Note that the calculated fission product distributions in the particles were only slightly different for the various conditions. Specific activities per gram of fallout vary directly with yield and indirectly with the amount of soil, but the relative fission product distribution among the different-sized particles is not greatly affected by these parameters. This reemphasizes the importance of particle size distributions and chemistry in determining the fate of fission products during falle * formation.

Table B. 5

VARIATION OF FISSION PRODUCT ABSORPTION WITH SIZE

AND HEIGHT OF DETONATION EFFECTS FOR A

100-KT FISSION YIELD DETONATION

	Cei	sium Absorbe	ed (g)	Zirconi	um Absorb	ed (g)
Particle Radius (cm)	Standard	1/100 Soil	10 MT, Total Yield	Standard	1/100 Soil	10 MT Total Yield
0. 0003 0. 003 0. 0154 0. 130	14.63 112.44 46.88 1.042	14. 76 113. 20 46. 06 0. 954	13.84 110.21 49.55 1.387	0. 444 33. 73 69. 28 10. 06	0. 539 35. 54 67. 76 9. 67	0.432 32.86 68.41 11.81

The computer program can also provide fission product profiles. Calculated profiles for both the 137 and 95 chains have been investigated for a 100-KT fission detonation, using 200° temperature increments and the particle size distribution given in Table B. I. The calculations were initiated at 2700°K for the 137 chain and at 3100°K for the 95 chain. The 900°K profiles for the 95 chain are presented as a function of particle size in Fig. B. 3, and those for the 137 chain are presented in Figs. B. 4, B. 5, and B.6. An unusual abscissa scale is used in these figures solely for the purpose of demonstrating the calculated values more efficiently.

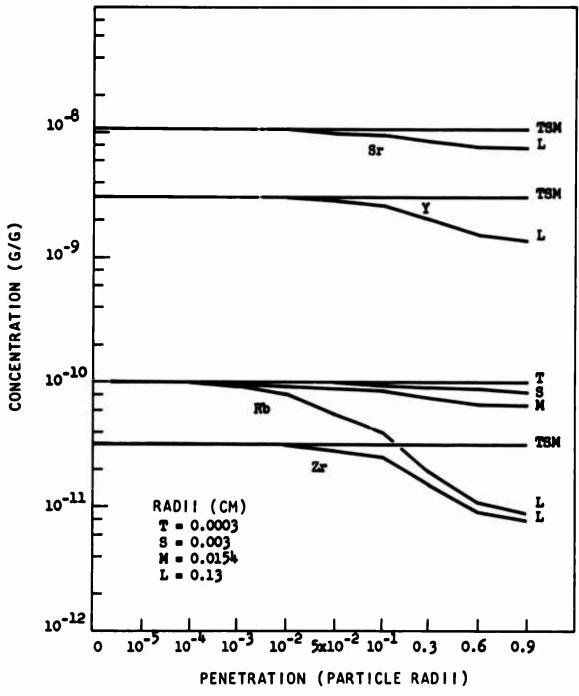


Fig. B.3--Concentration profiles at 900°K of 95 nuclide chain members as a function of particle size for standard detonation

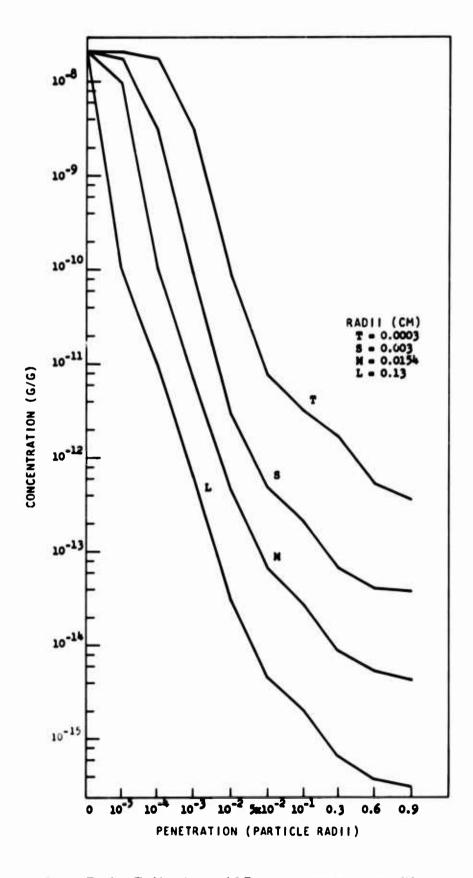


Fig. B. 4--Tellurium-137 concentration profiles at 900°K for standard particle size distribution and standard detonation

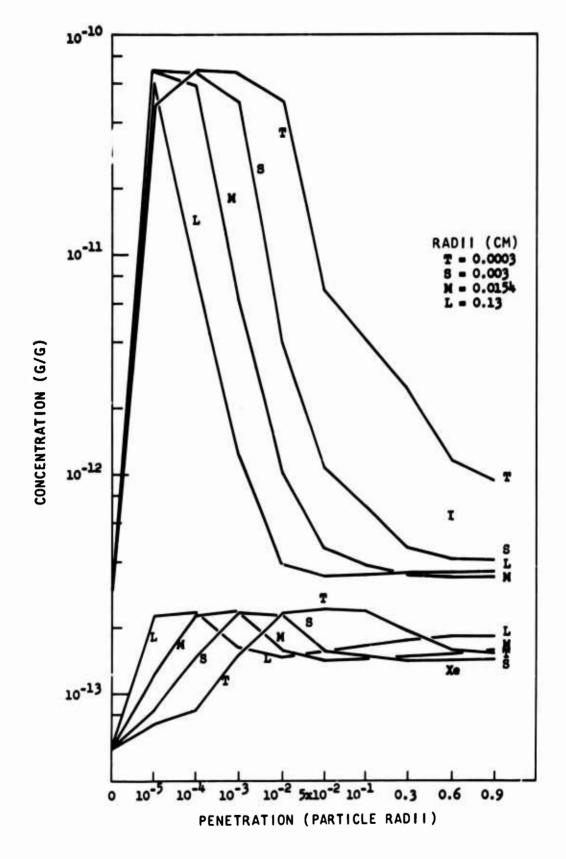


Fig. B. 5--Iodine and xenon-137 concentration profiles at 900°K for standard particle size distribution and standard detonation

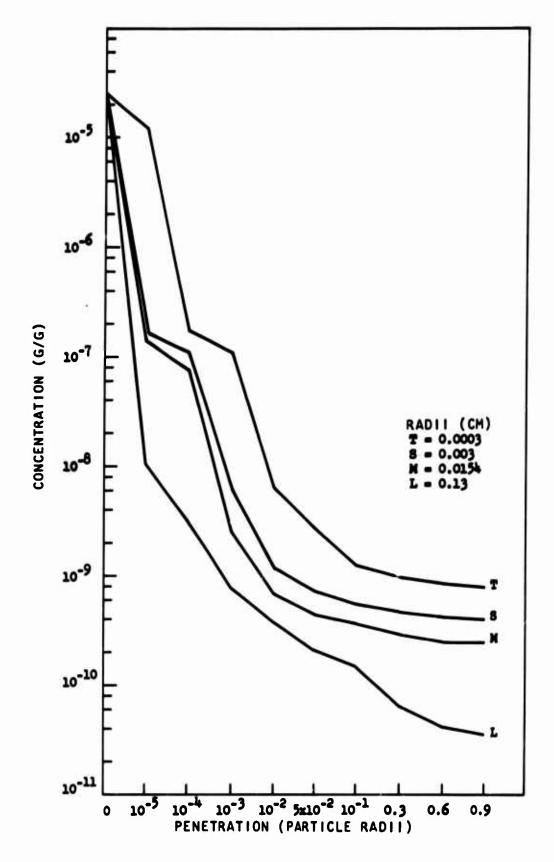


Fig. B. 6--Cesium-137 concentration profiles at 900°K for standard particle size and standard detonation

The 900°K profiles for the 95 chain are very simple. They show only small deviations from a flat profile for rubidium in the S, M, and L particles and for the other fission products just in the L particles. The 900°K profiles for the 137 chain are structurally more interesting. The distribution of tellurium within the standard particles at this temperature is shown in Fig. B.4. Calculated tellurium concentrations span eight orders of magnitude in the L particles and five orders of magnitude in the T particles. The concentrations have been drawn as a continuous curve in these diagrams as an aid to the reader and are not to be used in interpolating. The tellurium profiles, then, are characteristic of a moderately volatile fission product (1% condensed at 900°K in this calculation) with a moderately high Henry's law temperature coefficient.

The iodine and xenon profiles (shown in Fig. B. 5) are even more interesting, since they exhibit maxima and minima as a function of radius and particle size. The sharp maxima in the iodine curves can be attributed to absorbed tellurium decay and subsequent diffusion and vaporization of the resulting iodine. The maxima in the xenon curves can be attributed to absorbed iodine decay and diffusion and vaporization of resultant xenon.

The cesium profiles are shown in Fig. B. 6. This set of curves would appear to be exhibiting fine structure that is somewhat dependent on the calculational model. The near-surface behavior of calculated profiles would seem to have an extraneous structure that is a function of the size of the time-temperature steps. The use of shorter time steps should smooth out this profile considerably. However, the general structure of these profiles seems reasonable.

It is interesting to consider the penetration depths used in these mathematical profile calculations. The particle radius $0.0003 \, \mathrm{cm} \, \mathrm{times} \, 10^{-5} \, \mathrm{fractional}$ penetration gives $0.3 \, \mathrm{A}$, a value too low on the atomic-size scale to be a credible penetration depth. This can also be said for the $0.0003 \, \mathrm{cm}$ radius particles at $10^{-4} \, \mathrm{radius} \, \mathrm{penetration}$, and the $0.003 \, \mathrm{cm} \, \mathrm{radius}$ particles at $10^{-5} \, \mathrm{radius} \, \mathrm{penetration}$. A value between $15 \, \mathrm{and} \, 30 \, \mathrm{A}$, a

borderline penetration, is the next smallest penetration depth plotted. The subsequent minimum value is 130 Å, a value such that diffusion is reasonable to consider.

A plot of the iodine profiles for the 0.0003-cm radius particles, T, as a function of temperature is shown in Fig. B. 7. These curves demonstrate a general decrease in total iodine loading as the cloud cools. This can be attributed to the expansion of the cloud, high iodine diffusion coefficients, and a low iodine Henry's law temperature coefficient. At the lower temperatures, where iodine is not very mobile, the iodine profile reflects the parent tellurium profile. The difference between the 1100° K and the 900° K profiles is caused by additional tellurium decay to iodine.

This orientation investigation demonstrates the effects of condensed state diffusion in a fallout calculational model embodying this phenomenon. It is interesting to compare this system with the Miller model. where diffusion is deemed infinitely fast above 1673°K and negligibly slow below this temperature. For the case of the highly absorbed nuclide chain, the only noticeable difference is that the calculations for the larger particles in the diffusion-controlled model suggest that equilibrium conditions may not be attained by these particles during fallout formation. Equilibrium conditions are assumed by Miller. In the case of a volatile chain, the diffusion model predicts an absorption behavior which varies with particle size. Miller's model assumes no absorption difference with particle size. In the diffusion model, cooling rate will slightly affect absorption, while in Miller's model it will affect absorption only to the extent that this quantity interacts with fission product decay. Radial concentration gradients are predicted within a fallout particle with the diffusion model, while with Miller's model there are no concentration gradients within a particle, only at the surface of the particle. It seems reasonable to state that the disfusion model provides a more revealing and perhaps a more accurate calculational tool for considering fallout. It is, however, a limited model, and the assumptions underlying the development of the model should not be ignored.

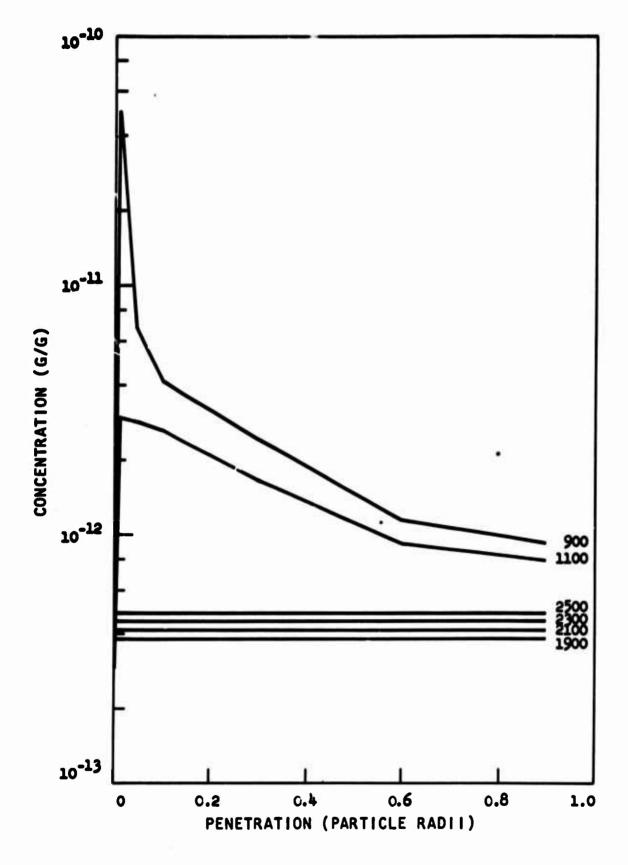


Fig. B.7--Iodine-137 concentration profiles as a function of final temperature for the 0.0003 cm size particles and the standard detonation

APPENDIX C CONTROL CARD FORMATS

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APPENDIX D SAMPLE OUTPUT AND SUBPROGRAM LISTINGS

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00122	.64	U	
00123	50.		DO 20 N=1.LC
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00220	98.		DO 401 K=1,1R
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00224	į		TO 401
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00227	į	100	SUM=SUM+UM(L.J.K)+#I(K)
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16200			SUBSECTION TO CALCULATE THE SURFACE CONCENTRATION (SURCON) FOR
10200		٠,	EACH 15010PE AT TIME STEP O'E.
003.46	102	,	Design III / John State of Company
00235	101		Do Sou Latell
00200	100		
00242	105		DV01640[(L)046/4[L-1]030K+kW(L-1)
59200	106.		65 10 502
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10272	3 2		TO THE OTHER PROPERTY OF THE OTHER PROPERTY OTHER P
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00272	137.	u	
00273	136.		CALL OCALC(1:K-1:1-E-300-DYEM)
00274	139.		KAIKZA+1
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200			CALL BEALC(1) KRIK. 1.1.41.Fig.DTER41)
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220	191.		DO 6013 ISU3=1.172TH
	152.		CONCILE 44 15031 = CONCILE 44 15031 + COVILE 44 + 717 + 504 CONCILE 4 (1-4.) + (1530 F11.) 215.3
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00322	157.		A-SUN
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00323	161.	v	
90325	162.		DO 603 HELEZFAIK
CX	2	۰	
		٠,	WALUE OF NEL 15 THE SUBSCRIPT OF F PREPRETSSON OF 12, 194E L.
2	100	,	MARKET THE PROPERTY OF THE PRO
	3		
00322	168.		NET = 1871 - 1
00332	169.	v	
00332	170.		
90332	171.		ייים בייים ביים בייים בי
00333	172.		DO S ILERETONI
00330	173.		M22=NTI-4E-II
	•		
00380			• • • • • • • • • • • • • • • • • • • •
00340	177		THE OFFERN GATH COMM.
00340	176.	1	AT NITHE.
00340	179.	u	
19600	180		WZII-PEL
90342	181.	The second second	1(2) 2-1
200	182.	101	IF (MEL-E9-2160 TO 40
00343	183.		N3EN2-1 (2)-1
			;
100	180	-	1 (CEL-15 0 10 40
00352	1		1(8)3-1
00353	188	121	IF (NEL. E0.4) GO TO 40
00355	189.		N3114-1(+)-1

00362 195. 00362 195. 00362 195. 00363 196. 00364 201. 00366 201. 00366 201. 00366 201.		
		IF (MEL. E.B. S) 60 TO 40
		1-(6) 1-1
		DELTA FINDS THE DECAY TERM ASSOCIATED WITH THIS PATH.
		The same of the same of the same of
	-	בשרה תברועו הברועו ושביווים להיו
		BCALC FINUS THE DIFFUSION TERMS ASSOCIATED WITH THIS DATH.
		CALL UCALCINELINTINE, I. 1.80F.000.9TERM)
		CALL BCALC (RELINITATIO) BOFI, FW, DYEAT)
		IF(NTIME.ME.NTOT)GO TO 19
	1	THE THAT I WE THE CONCENTRATION AND PROFILE POINT
L		· CALLEGE LOCAL
	ı	DO 14 VELVIE
00373 209.		IF(IDPTH-E0.0)50 TO 1A
		DO 160 ISUBELLIDETH
		CONC (L.M. 15UB) = COMC (L.M. 15U3) +POFLT+SURCELINEP. N22) + (PROFIL (UFFR)
00400 212.		-R(M)-6(150B))-PROFILIDYERAI-R(W)-5(150B))
	-	CONTINUE
		16 AVEUITLE AT EAVEOUTLE STEET SCHOOL PERSON SEPTINGS FOR THE TOTAL
		8J6=5JF-3JF1
		B IS THE SUM OF ALL CONTRIBUTIONS FOR THESE STATES.
00400		- McG-90-90-90-90-90-90-90-90-90-90-90-90-90-
DAMAN SSA		
- C	:	TO 10 10 to
		Tight
00014 223	•	1(5)=(5)+1
1	1	TOTAL STATE OF THE
		A
		15(1(6)=44)131-131-131-131-131-131-131-131-131-131
	=	1(3)=1(3)+1
ĺ		IF(I(3)-H3)1116111610
	. 10	1(2)=1(2)+1
l.		IF(1(2)-42)101-101-10
	16	6=B+SURCO.4(MEP+M22)
Ĺ		E+INSENDS
	. 962	IF (ALDGEOTHENIL MIT ET) +ALDGEOTORY) . 67 . 75.1 GU TO SUS
1		DVD=DKK *# XOL (L) *HEN (L, *IT ME) **HIL. NIT (E)
		504 01 05 04
_		DVD=1.E57+E4(L.NTINE)
.007	*	CONTINUE
		IF (3V)) 666+ 1001-666
	§	TYTELLO (LONT IAE) —A
102 56400		TESUC-MILE.
242		SURCON(L.MTIME) = TYYOUD
AMARY SEE		Tributanting of the state of th

00065 2 00067 2 00077 2 00074 2 00074 2 00074 2	2.6		IF(IDPTM.Eg.0)50 TO 47
			DO 470 ISUBEL IDPTH
	247.		CONCIL. H. ISUB) =CONCIL. M. ISU9) +SURCO: (IL. MTOT) +PROFIL (BFC (L. MTOT) +T(
	240	47.	ONIOTOR(M) DO ISUSO)
	250		47 AVCON(L. 4) ZAVCON(L. 3) 4SIBPONIT - 21701) EIPIT - 21701 -
	251.	900	CONTINUE
	252.	v	
	255	ی	PRINTOUTS
	255	ט נ	
I.	256	þ	SUBFACE CONCENTOR DOINT ON THE
	257.		
	256.		WATTE (6, 1075): KSIONCSZO (MAME (L.) OLEJOLC)
	259.	1	DO 1090 J=1.NTOT
2 21500		8	BATTE (6.1076) TEMP (J) . ATM (J) . (SUBCON (L.J) . LETLC)
	26.0	,	TO USE A LICENSE
ñ.	263.	·u	OF THE STATE OF TH
TV.	. 192		WRITE(601200) MCS10) MCS20 () (AME (1.) (= 10.10)
00533 2	265.		Do do Jalintor
	266.	1	<u>00 81 CST/IC</u>
	.67.	10	6AS(L)=SURCON(L,J)+FE4(L,J)+VOL(J)+H40L(L)/(RCST+TE4P(J))
	266.	90	exite(6.1076) Tear(J) - ATV(J) - (GASTL) - (C)
	269.	u	
2000	27.5		AVENAGE CONCENTRATION PRINTOUT.
		اد	
	273.		DO 43 JET 18
	276.	83	#RITE(6.1202)R(J). (AVCONIL.)J.(E1.1C)
	275.	v	
00570	276		LOOP BE FINDS SURTARY RESULTS.
	27.8	,	
	279.		SOLID(L)=0.
l	-092		DO 89 K=101K
	281.	69	SOLID(L)=SOLID(L)+AVCON(L,K)+HT(K)
	282		AVAVIL) = SOLIDILI) / THE
	265		10TAL(L)=50L10(L)+645(L)
00615	204	0	01F(L)=(01AL(L)=(101AL))
n e			WALLE OF ALL AND AND AND AND AND AND AND AND AND AND
	267.		MATTE(6,1211) WCS1-MCS2-TEMP(NTOT)
	-992		
347	209.	U	LOOP 250 PRINTS PROFILE POINTS.
	-063	5	
00632	31.		DO 250 L31,LC
	202		DO 251 KELLO LELEVINARE (E. O. 101KF) OF SELEVIDATED
I		162	WELLE GO S THE KI O CONC. I SEE THE PERSON
2 25900	3	8	CONTINUE
Γ	2	1	
96657 2	22	J	PAINT SUMMARY RESULTS.

19900	299.	252 MRITE(6-1204) TEMP(NTOT) -NCS1-NCS2- (MAMEIL) -L=1-LC)
00672	300	
00200	301.	#RITE(6,1206)(SOLID(L),L=1,LC)
90400	302.	#RITE(6,1207) (TOTAL(L).L=1.LC)
00714	303.	WRITE(6,1208) (YLD(L, HTOT), L=1, LC)
00722	304.	##ITE(6,1209)(DIF(L),L=1,LC)
00730	305.	IF(IL6.E9.0)60 TO 3
00732	306.	CALL FINAL(AVCON-BAS-R-UT-NAME(1))
00733	307.	60 TO 3
00733	304.	
00733	309.	C EXRON STOPS
200	310	
200	3111	1000 GAT 15 (0-1050) NAME (1, 1-4-0) (1, 1-1) (1-1)
22.00	312	CALL DOWN
98734	-	X Z Z Z Z Z Z Z Z Z Z Z Z Z Z Z Z Z Z Z
00757	315	CALL DUMP
00757	316.	
00757	317.	C FORMATS
75700	316.	-
09200	319.	1050 FORMATITHO.SK.35H ERROR. PROGRAW CAWAT COUTTIAN -/104-29HTIE SU
99,00	320.	XRFACE CONCENTRATION OF ANIMALISMIS INFINITE./ISX.154-NOLECULAR WEI
90760	321.	AGMT = .E16.8/15x.23+4E'RY'S LAW CONSTA'IT = .E16.8/15x.5H4" = .E10.
09/00	322.	X8/15X-17HTXE TNCHE-EIIT = .E16.8/15X-14HTEMPERATURE = .F9.1/15X-11
09200	323.	XHVOLUME OF CLOUD = .E16.8)
19400	326.	1075 FORVATITHISOX SINSUNFACE CONCENTRATIONS IN UNG 115X MICASE 15.
00761	345.	
29,00	326.	1076 FORMALIANTE STATEMENT OF THE STATEM
2	361	STATE OF THE PROPERTY OF THE P
2760	920	
00764	330.	1261 FORWATTHIS SAFEANERING CONCENTRATIONS PER PARTICLE SIZES INTO THE
00764	221.	.ASE NJ 13.1M A4/49X.21MAT FIMAL TEMPERATIME . FF. 3.11M DEG KELVI
00764	332	.N/728K+7HISDTOPE//2×-15HPARTICLE RADIUS/6X+7+(14 C4)+13K+A6+5(13K+
900764	333	.Ab) I
00765	.34.	FORWATTIN SERVES SESTING SESTI
99200	335.	1204 FORMATCHILITY STATEMENT OF RESULTS AT FINAL TEXT OF STATEMENT
90200	.336	.ELVI I DX 94CASE 40. 13.14 . A47/ZEAX /HISDIOPL/ZEAX AGESTISA PAR
00767	337.	
00770	•	1206 FORMAL CHOINTENANT IN SOLID PHASE INTELLEGISTICS IN THE CONTROL OF THE CONTR
		TOTAL THREE TRANSPORTER TO SECURITION OF THE SECURITION OF THE SECURITION OF THE SECURITION OF THE SECURITION OF THE SECURITIES OF THE SECURITION OF THE SECURITIES OF THE SEC
2773		FORMAT (1HO - 10HD IFFERENCE - 7X -
512	3	1211 FORMAT (IMISEX) 264P R O F I L E P O I M T S-21X-94CASE 110. +13+13
00775	344.	AT/750X: 32HCONCESTRATIONS AT VARIOUS DEPTHS/49X:21HAT FINAL TEM
ST13	345	
00776	200	IZIZ FURMATTIMOPDEN SHROOPADFIDM ISSIONE
44400	346.	1213 FORWATTIX-OPF9.6-1X-11(1PE11.5)7
	5	

	INITAL													
DI FOR INITAL COMPLETION BY UNIVAC ITOE THIS COMPLETION MAS DONE	ATION K		LIDE FORTRAN-IV DAT	C 66 AT 2	1:06:04	HOV 1966 F4006	F4098							
SUBROUTINE INITAL	NE INIT		ENTRY POINT 001103	001100										
STORAGE	USED (B	LOCK. N	STORAGE USED (BLOCK, NAME, LENGTH)	(H)										
83	1000	CODE D	000626											
88	0002 *8		000023											
EXTERNAL	REFERENCES		(BLOCK NAME)	6										
8	9000	ALOG												
88		OIFFUS EXP												
28		NEDUS NECUS								4				
38	H	MIC25												
8		REAPS												
STORAGE	ASSIGNA	ENT FOR	STORAGE ASSIGNMENT FOR VARIABLES TBLOCK	S TBLOCK	٠	TYPE RELATIVE LOCATION HAZE	LOCATTO	We HAZED						
1000	000054 10L	104	0000	000164	1 Inde	0000	0000273	3 101F	6000	012000	1201	600	6.6347	4
	191000	104	0000	000505	. L.	0000	00500u		9309	UN0534		2000	D. 6541	4
ı	000545	1114	0000	00056.3	_	1000	400000	1.5	0001	009703		1000	0.0507	1.0
1000	000000	9041	1900	40000	9951 /	1000	000625		1000	401-10		1001	n n n n n n	ž.
9900		2002	1000	201000		2000	77000	1 20r1F	0000	306,	Zi-u-Zi-	1939	2000	20.03
1000		5606	1000	000350	_	1000	000473		1000	093507	300	1007	A.05.32	36.36
1000	000554	3676	1000	000612	2 4046	1000	004635		1000	043716	1	100'	ACT 0.72	1235
1000		5106	1999	00100		1000	100000	9 1 2	1000	916160	2	1664	2000	
1000		759	1000	0000		1000	000542			4 901334		4900	4 21 00003	100
2000	001626	ATA C3	0000	R 000131	T. C.	0000	R 000160		1	4		3000	H 054124	5
B 0000	201000	400				- 1	TO DOUBLE	D DIECK	0000	4 000 3 322	100	1000	020000	
	900000	¥			100				2000	2 00 3 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5		1090		
H 0000	000134 HTM	H.	0000			2000	00000		0000	9000-0-1	11.5	-1362	THERMET	0
1 0000		,	-1	21000 1		0000	1 999156	- [0000	040133	-1	1.001	I POODO	۲
0003 I	-	TOTA			y Ld 9	0000	2 00000 K	Par Par	0000	300010	25.7	1000	611926 I	
		-:		R 001556	2000	0000	R 000152	1		1		1000	ntie151	Trail
- 1	1000		- 1				- 1	-1	- 1			2002	P ncons	
2000 K	961100	8	2000	N DOUGED	S STAGE	9560	R DETENDE	18.0	6000	2.000000	-	-	-	-

					-	
0000 K	R 000000	TLU 0000	R 000115 2DC	0000 R 99014c 207	11052 TALONG & F-000	11: 401030 & 0000
10100	:	SJBROUTINE I				
00103		COMMON YLD(6,40), MAYE	0.401.34 E(6).445 (6).DFC	6) - 440L (6) - DFC (6-40) - HEN (6-40) - AK (6) - K (30) - 4	K(6),4(30),4	
00104	•	COMMON/SIME/LC+IR+NYOT		NCS1 NCS2 TYT ILG. TOP FH. D		
90100	•	DINENSTON OCTIV		The state of the s		
00110		DATA CI74.66E37.C472	E37.04/2.3416261E-	#16261E-5/1-1/3-14159265/		-
01100		AT IONAGOE	1	CALCULATES THITTAL TOPOLS	S AREJEJ.	
0110	:	VAKIABLE	DEFINITIONS			
01100	13.	3	LEMETH OF CHAIN.		Ď	
00110	14.	TKI	GF DEVIC	USED IN CALCULATING YELLUS		
00110		GAT	ED FOR	CALCULATINS VOLUME, L'ISHF.	, TEMPERATURES,	
01100	17:0	1	MUMBER OF PARTICLES	(IN KT)		
01100		NAME	100	NES OF		
00110	19. C	AirOL.	COUTAI	3		
00110	20.	r.	PER 10:000 FISSIONS	TIAL YIELDS ANT TEU.	410 15	
01100	25.	OFC.	N GF DI	FUSION COEFFICENTS.		
200	38	HEN.	ARRAY OF HENRY'S	ENRY'S LAW CONSTANTS.		
00110	55.	í	1/2)=Ke	LIFE)		
01100	299.		OF PARTICLE	E		
			2	HIS		
	29	₫	ARRAY OF CLOUD VOLUTES	INCREMENTS. (TH SEC)		
01100	30.0	ES/B	Cr TENPE	Ē		
00110	31.	£	CF HALF-LI	(IN SECONDS)		
	35.	Total	TOTAL MINHER OF T	TIME TEADEDATION CLOSE	ICLE SIZE.	
011100	. K	HTNP	TARTING	IT TEMPERATUR		
00110	35.	CTMP.	L (LOWES	TEVE		
001100		Tage	TOTAL SETSAT OF PA	TEMPERATURE.		
01100	36.	ATA	Ę	DAE OF CLOUD.		
		TALIFIE	A"CKATS OF COUSTAIN	CONSTRAINTS USED IN CALCULATING HENRY'S LAS	HKY'S LAS	
00110		001,002	ARRAYS OF CONSTAIN	CONSTAINTS USED 1" CALCULATING OF	oteFusta	
1100	, 45.		COEFFICIENTS. COS	COSTOPETI = DC27TEMP.		
00 10	T	READ THOUSAND AND BE	A TA A STATE OF THE PERSON	THE TOTAL BUILDING TORE		
001100	*5. C		100 min on unit	A C. PARA T. ICRO.		
1100		READIS-115) (NTTTLETU)	INTITLE (J) . J=1.12) . HCS1.NCS2	KS1.NCS2		•

90136	.6		READ(5+2002)(Q(1)+1=1+10PTH)
94100	20.		10 WATTE(6:100) (HTTLE(J) , J=1:12) -HCS1 -NCS2-LC+TAT+34T
00157			00 40 Later Color (1) 202 (1)
29100	ž		READ OF COULTING CALL CONTRACT
20100	ŀ	ſ	08_ <u>=511E(8+191)MV4E(F)+#MOF(F)+AI(F)+HC1(F)+HC5(F)+D</u> C1(F)+DC1(F)+DC3(F)+HF(F
	55	١	
1207	8		READIS-2002)HTWP-ET-P-TEMPV
1214	57.		READ(5,2003) (R(K),PX(K),R=1,1R)
00223	28.		READ(5-2004)1#67-11V
1227	29.		COM 0464 (UKT) 44 (575)
1227	3	u,	
00227	.19	u	ŀ
00227			ACCOUNT SINGLE S
00227	3	u	INCREMENT
00227	:	v	
0230	65.		22 ZTMP=i(TMP-TEMPV/Z.
16200	.99		21744L06(112474) PS
0232	67.		DS 20 L=1.LC
0235	-99		20C=10.ee(DC1(L)=DC2(L)/Z14P)
0236	.69		ZDT=ZDC+ZTM .
0237	70.		ZSOXE0.
00240	:	0.5	20 ZSUM=ZSUM+DIFFUS(ZDT+R(1))
2920	12.		IF (485 (250M=C) : 61 - 7 - E - 61 62 13 21
00244	ż		H TWP=HTMTEMPV
0245	:	•	60 10 22
00245	2	u	25-14-1016 TO 101 MM 101 TO 101 WAS ASSESSED.
00245	ė	u i	PKINI METONI PENTENDER NO VOCORE TANATALENS.
00245	:	u	
00246	2		21 BKI E (00102) HI
00253	2		107 - 170 -
6520	3	v	THE OWN STATE STATE OF STREET
0253	.19	٥	1
07253	95.	u	
9520			141-101-101
/520/	:		
00262	95.		91 HI(K)=TBI+PH(K)
10264	-98		MAINECOSTUATION
00267			
21200	98		STIFF (PITOS) KIKI - ALIKA
00277	8		170-11-11
7200		,,	COMPUTE INITIAL TIME (TINT). DICREMENTAL TIMES (T(J)). AND
11200		4	TOTAL TIMES ATMICUST.
00277	3	·	
1		1	TIMETALOGRAPHY/CI+C21117C3
10000	ģ		ij
200.00	8		
00303	6		64 TIENP=TEMPL-TEMPV
200.00	-		1
20200	g		T(J)=ALOG(TTEMP/TEMPL)/C3
00306	100		IF CT TEMP-ETMP161-60
17:00			

	103.	62 TEXPLETTEMP	
00316	104.		
ŀ	105.	63 WRITE(6.106) TE-P(40)	(04)
00322	100.	NTOT=40	
1	107.		
	108.		
	109.	05 ATW(1)=TINT+T(1)/2.	7.
	110.	DO 66 JEZINTOT	
1	.111	66 ATM(J)=AT4(J-1)+	AT#(J)=ATA(J-1)+(T(J)+T(J-1))/2.
		5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	
1	113.	PRINT TIME-TEMPE	PRINT TIME-TEMPERATURE DISTRIBUTION.
		1	
	115.	#ATTE(6,107)14TOT	
	•	שלייים יים יים יים יים	
	-	70 8411E(6+108)AT*(J) • I(J) • TEMP(J)
		SKITE (6-114) TIT	
		The state of the s	
		CONTOIL MILITAL	CONTOUR MILITAL MELLOS IN GRASS.
		41.15	
	166.		
	TO THOUSE	1	THE TEST FOR COURT IN THE EN WHAT THE COURT WOUND SE VENITHERING.
		1F (ML(L) : 0.0.0	
1	126	AK (L)=.69313/HL(L)	
1	.034	-1	
70500		70 TICLICATION (L) OTKI OTICL	• K.• T.C.
	1	- 1	
	130.	CALCIA ATE VISING	C al craot true.
L	135.	1	:
		H.M=1.E38	
00366	133.	DO IN LEISLE	
	134.	IF (4_(L) .LT.1.E-5)	5) 50 T0 14
	135.	HENSAINI CHEN. METERS	(L)
I	136.	14 CONTINUE	
	137.		
1	130.	IF (HLN.GT.TINT)	H_WITINI
10000		2000	
	181	0-70L	
00410	163.		
		16 IF (LC. FO. 1) 60 TO 10	
	145.		
	146.	73 YICU =YICU+0	=YI(L)+30P(L)+YI(L+1)+(1,3P(L+1))
02400	147.	18 TILLCIETICIADDPILCI	butth
	146.	TT=TT+HCH	
	149.	IFITICATION 69 TO 16	0 TO 16
	150.	IF(KRUZ.Eu.1) 60 TO 13	70 13
00426	151.	IF (ABS(TT-TIMT).LT.1.2-5)	LT.1.2-57 69 T0 13
00000	152.	KRUZ=1	
10000	155.	HCMST INT-TT+HCM	
Ĵ		67 01 00	

96 400	150.	17 YLD(L+1)=YI(L)	
00436	157.		
90436	156.	C PRINT OUT ADJUSTED YIELDS.	5.
90430	129.		
24400	1001	A WKITE(De 109)	
54400	162.	75 WRITE(6:110)HAME(L) - YLD(L-1)	
54400	165.		
00445		C CALCULATE DIFFUSION COFFFICIENTS	FICIENTS AND FIRM'S LA CONSTAINS.
00452	100	Do Su Jetanfor	
95400	167.	1	
09400	100.		2(L)/TE*P(J))
19400	169.	60 DFC(L.J)=10.00(DCITE)=0C2(L)/TE4P(J))	2(C)7fE4P(J)1
19400	170.	C PAINT DIFFUSION AND NE'N	Y'S LAS COEFFICTEITS.
19400	122		
00064	K	PATTERNAL TO NOT COMPANY OF THE PARTY OF THE	
00474		30 50 J=1+HTcT	
00477	175.		ערים) ירבו ירכו
90500	176.	50 mRITE(6,113) (HEN(L,J),L=1,LC)	11-(3)
90200	::		
00500		ביירות אור אחרתה בתי ארר וו בזי	11 630
51:00	100.	CONT. 4564EBBPT + BKT	
0.590	181.	CONI=.176((SKT) ee(=.375))	
.00217	185.	DO 40 JELVINOT	
00522	183.	40 VOLTJECOTOTEXPTCONTOATH(J) 1) 003	2000
92500	100	2000 FURNATION 51	- 1
50000		TO A II I TO POST TO THE POST OF THE POST	//••• V T V O
00525	187	- 130 10X 150 10X 150 10X 1	APP NO WISHINGTON TO THE STATE OF THE PARTY
00525	198.	/64CFOR TIME-TEMP DISTAND	/64(Fox 114E-154P)157819011911//56xe19415070P. 0402 4 FF35//47 64413
90525	169.	THE . BX . 9H YOU ECULAR . TX . THI	746.9X • 9443LECOL 48 • 74 7 7 1 1 1 1 1 1 1 2 4 • 3 1 1 1 2 4 • 3 1 1 1 2 4 5 1 2 4 5 1 2 4 5 1 1 2 4 5 1 1
00525	190.	/DC2+12X+4mHALF/22X+6HIEI	/DC2.12X. 4 MALF/22X. CHIETGHT. 94. SHYTELD. TOX. 4 HLTEL/2UX. 941 G/ NET.
00525	191	VKE TSHEATURE/ITEDROPESSY	74.6.1.36.1.41.345.1.1.6010.63X.9.4(SEC0105).735T.94+TSST045171
9750	192.	2001 FOR 4AT (A6. 7F13.0)	
75500	2.5	101 FORMAT (SAMAGE 12X - F4. U. 3X	101 FORMAI COACAGO LEXOPE-US SXOPT. IS GX OF P. SOTY OF G. IS TASTR. COTAGE TO TAST
00230	6	Z002 F384ATTRF16.01	***************************************
00531	196.		6 TEMPERATURE = "FR. 0.12H DEG. KELVILLINK.
16500	197.	ZUHPINAL	0.00 HZH CES.
00531	196.	/CREMENT = "FO.O.12H DEG.	KELVIN
00532	199	2503 FORWATCZFIG.0)	
00533	500.	2004 FURMAT(2E16.6)	
90534	201.	103 FORMATCIOX, SCHITCTAL PART	FORMATCHOK, SKITOTAL PARTICLE ACTOM = , IPC16, 0,64 5-4/47
00535	203	TOT TOWNS THE TOTAL PRINCIPLE TRANSPORT TO THE TOTAL PRINCIPLE TO THE TOTAL PRINCIPLE TRANSPORT TO	TATABLICAS INC. 120 IX OLIMPARTICLE SIZES/
00536	204.	1u5 FORWAT(10x,610,6,5x,13716,2)	(-a)
00537	205	106 FORWATTTHOUSE, SOUTHERS	NE TOO VALVE TI WESTERNOR STATE OF THE PARTY FOR
00537	206.	/ED./10X.40HTHE LOHEST	/ED-/10X+40HTHE LOHEST TEMPERATURECONSIDERED ALL CESTAGES SPECIFIED
00537	207.	/GREES KELVIN	
00540	208.	107 FORMAT (1HU. 91, 29HTI-E-TE	FORMATITHUS 93, 294TI-E-TEMPERATURE DISTRIBUTION, 1X, 12, 74 STEPS, 1/12
MACE	2000	THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, T	THE PERSON NAMED IN COLUMN TWO IS NOT THE OWNER, WHEN THE PERSON NAMED IN COLUMN TWO IS NOT THE PERSON NAMED IN THE PERSON NAMED IN THE PERSON NAMED IN THE PERSON NAMED IN THE PERSON N

	210.	LUG FURNATULM .9X. FB.S.G.K.FB.S.10X.F7.11
00542	211.	109 FURNAT (1H0.20x.41HYIELDS: IN GRAIS ADJUSTED TO START TIME //10x.7H
00542	212.	/ISOTOPE:10x.5HYIELD)
00543	10543 213.	110 FORMAI (1H , 6X, 46, 6X, 17E16, B)
99500	214.	111 FORMAL (1H1,24X70H)]FFUSION COEFFICIENTS (LTHE CHE) ALC HENRY'S LAW
99500	0544 215.	CONSTANTS (LINE TWO) - 3X - SHCASE NO 13-111 - 6A4//4X - 4H (E-140-2) - 5(1.54
44500	216.	(146)/)
00545	217.	112 FOR 4A1 (1H0-1X-F9-1-7X-6(3X-1PE16-2))
90500	210.	113 FORMAT(17X,6(3X,1PE10,2))
00547	219.	114 FORMAI (1MD+9X+56H5TART TIME FOR CALCULATION OF SUPERALE CONCENTRATI
00547	220.	/0 .5 IS .F3.44H SEC)
00550	221.	115 FURNA: (12A6-13-1X-A4)
15500	222.	RETUKI
00552	223.	Eid

CANTING FINAL CHIRK PORTARI-2V JATED 11 NOV 1966 FROME	OI FOR	FINAL ION BY UITU	AC 1108 I S DONE OF	ORTRAN-	TE L		Choo							
VARIABLES (3LOCK, TYPE, HELATIVE LOCATION, HAVE) VARI	THIS COM					00.	0116							
RAME EMINAL EMI									,					Ť
0000 1000	SUBRO	UTINE FINAL	ENT	RY POINT										
STATE CONTRIBUTE CONTRIBU	STORAL	JE) 0356 39		E. LENGT	(+)		i							
UNIVERSENCES GLUCK NAME				152										
ERNAL REFERENCES (BLUCK: NAME)				900										
### ##################################	EXTERI			CK. NAME			1		•			9		1
NAME ASSIGNMENT FOR VARIABLES (30,00K, TYPE, MELATIVE LOCATION, MAYE) 130,000			35 15											
100 100.234 100F 100.00 100.0	STORA	1 1	9. A	AR TABLES	(BLOCK.	1	LOCATION							
100	0000	000234	00F	0000	20		410000	1076	1000	210015				9
100 1000227 K	0000	A 000074	277		00000000000000000000000000000000000000	0000		1546 1	0000 H	099909		۲ -		
1. SUBROUTINE FLIGE(TAVCOTTERASTATERINE) 2. DINELSION AVCON(6.34).665(6).4(30).47(30) 3. DINELSION PUTOTTERASTATIONS 4. DIVERSION FUTOTON (6.34).655(6).4(30).47(30) 5. DO 10 JUTIN FUTOTON (6.34).655(6).4(30).47(30) 6. COMPONING FUTOTON (6.34).652(1.74) 6. DO 11 KEINLC 6. DO 11 KEINLC 6. DO 11 KEINLC 7. DO 11 KEINLC 7. DO 11 KEINLC 12. DO 21 KEINLC 13. DO 20 LEINLC 14. DO 30 KEINR 15. DO 30 KEINR 16. DO 30 KEINR 17. DO 30 KEINR 18. DO 30 KEINR 18. DO 30 WEINR 18. DO 30 WEINR 19. TANGEN 19. FUTOTON (6.40) 19. VITO (6.40) 19. FUTOTON (6.40) 19. DO 40 I=1/R	0000	1 300227	TOT		0300351	0000			0000	30003		 P-		, i
1. SUBROUTING FINALTANCOTIGAS: Graterinage) 2. DINGHASION AVCONGES, Grates of STATES 3. DINGHASION PCTANTIFESTANTIATION STATES 4. CONTACT OF CASTANTIFESTANTIATION OF THE STATES 5. DO TO JEINT CONTACT OF CASTANTIATION OF THE STATES 6. IS ECS(J) = FCS(J) + AVCONGES, J) 7. DO TO TO FCS(J) + AVCONGES, J) 8. IS CASTANTIATION OF THE STATES 10. DO TO TESTANTIATION OF THE STATES 11. DO TO TESTANTIATION OF THE STATES 12. ZO GT=GT+GAS(L) 13. V=0. SO K=1.IR 15. V(K) = L(K) + AT(K) 16. DO TO TESTANTIATION OF THE STATES 16. DO TO TESTANTIATION OF THE STATES 16. DO TO TESTANTIATION OF THE STATES 16. DO TO TESTANTIATION OF THE STATES 16. DO TO TESTANTIATION OF THE STATES 16. DO TO TESTANTIATION OF THE STATES 16. DO TO TESTANTIATION OF THE STATES 16. DO TO TESTANTIATION OF THE STATES 17. SO VERVEYCK)		•												
2. DINEHSION AVCON(6.30).645(6).4(30).3T(30) 3. DINEHSION FCST30).FYISOJ.ZSOJ.VY30) 4. COMPON/SINGLESTRATOOT.NCSI.NCS2.TST 5. DO 10 J=1.TR 6. LO 11 K=1.LC 7. DO 11 K=1.LC 8. LO 12 K=1.LC 10. DO 20 L=1.LC 11. DO 20 L=1.LC 12. ZO GT=67.65(L) 13. VY=0. 14. DO 30 K=1.TR 15. Z(R)=K(T)=R(K) 16. DO 40 I=1.TR 17. SO V=VV+V(K) 18. DO 40 I=1.TR	10100	;	SUBROU	TWE FIN	IRLT AVCOTIBASE	- TO STOURKE)				51 140615				
3. CUM 40.4/SIGALC.IR. VIO.10.1/SIO.	00103	2.	DINCHS	TON AVCO		5) -K(36) - TL	30)			FI.139230				
5.	90100		COM 401	104 FC(2)	00101010101010101010101010101010101010	1 (30) • Z (30) • 1 • 12CS2 • Ts T	(05)			F1.13021C		ì		
7. 03 11 K=1.LC 4. 11 FCS(J)=FCS(J)+AVCON(K.J) 9. 10 C3MTiAUE 10. 6T=0. 11. 03 20 L=1.LC 13. 20 6T=6T+6AS(L) 14. 03 30 W=1.IR 15. 2(K)=K(T)7R(K) 16. 30 W=9V+V(K) 16. 03 40 T=1.IR	00106	in a	FCS (1)	JETO JR						T-005.35		1		
6. 11 FCS(J)=FCS(J)+AVCON(K.J) 9. 10 CONTINUE 10. 6T=0. 11. U3 C L=1.LC 12. 20 6T=6T+6AS(L) 13. 20 6T=6T+6AS(L) 14. D3 30 K=1.1R 15. 2(R)=K(I)7R(R) 16. V(K)=Z(K)+BT(K) 17. 30 V/=VYV(K) 18. D3 40 I=1.IR 18. D3 40 I=1.R	21100		1	CELEC						FI.146722			1	
10. 67=0. 11. DJ 20 L=1.LC 12. 20 GT=6T+GAS(L) 13. VJ=0. 14. DJ 30 K=1.IR 15. V(K)=Z(K)+Y(K) 15. V(K)=Z(K)+Y(K) 16. DJ 40 I=1.IR 16. DJ 40 I=1.IR	00115			=FCS(J)+	AVCONTACT				1	FI-440530				
11.	00121			ğ								,	8	
13. VFED. 14. DJ 30 K=1.IR 15. ZIKJ=KI17/RIKJ 16. JO V=VV+VK) 16. DJ 40 I=1.IR	00122			2701=7						FI.3437.50				,
14.	00127		V7=0.	2150						FI. A05 37				
15.	00130	14.	00 00	K=1.IR			•			FI IANSO				
17. 30 V/=VV+V(K) 18. DO 40 I=1.IR	00134	15.	7=(X)A A(K)=7	(X) *±T(K)	C				!		-			
	00135		1	VIKO				Western a warmen of the second		FI.4052U				ij
	00142	.61	FC 173							F1.1496.30				

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2770	•		FI KUDDO
0143	21.	3	FI IADZJU
2 4 5 0	*		
2**	. 77	C FAIN DAIA.	FI 140710
0143	23.		JC7041.13
0145	24.	#XITE(6,100,000 to 2001,000 to 2000,000 to	FILLED 730
0153	25.	D3 50 K=1+13	F1 14374C
00156	20.	50 EATTR (5-101) X (A) - FCS(A) - FC(A) - FC(A)	FI 140751.
0165	27.	RETURIA	DC-UTI IS
0106	20.	100 FORWAT(IMI-47X-22-LOUG-TERM RESULTS FOR -44-6H ISOTOPE-161-54-64-ELLAIGH	MICASE FILETON
0166	29.	-NJ. +13-14-+A4//9x+15-4PARTICLE PADIUS-STV-21-AVERAGE CO. CELTER TIOLS FIRST INTERIOR	(+T10.:0F1'141'11
0166	30.	-5X*27HTOTAL AVERAGE CONCENTRATIONS-SXSISHTOTS AND NO STATE AVERAGE CONCENTRATIONS-SXSISHED	4-1(C.1)FI (A102)
0100	31.	** 12X+20HDUE TO SOLIS'S DECAY SAF 20H (I TCL 19FS BAS TIBS) SAK	Selvision S
0166	32.	•2H IN PARTICLE/36X•7H(G4/54)•22X•7H(G4/G4)•20X•4H(G4)	- VI VI L
10167	35.	101 FURNAT (15x+17-5+12x+125)4-3+15x+12=14-3+15x+12=14-6)	100 C 100 C
0110	ż	CAL	FI .420.00

PILAT	TON BY	IL DAVI	106 FORTRAN-	JAC	E) 11	NOV 1956 FROM	FEDDA			-				ì
S COM	PILATION	MAS DO	THIS COMPILATION JAS DONE ON 19 DEC 66	¥	21:06:12						5			1
SUBRO	SUBROUTINE SCALC	ורכ	ENIRY POINT 00015	1000157										
STORA	STORAGE USED	'SLOCK.	MANE. LENGTH)	E										
		300	000225											
		PUATA	000030											
	2000	STANK	979100											
		BI INLOOP	023160 000001											
EXTERI	NAL REFE	RENCES	EXTERNAL REFERENCES (BLUCK! NAME)	6										=
	9000	DIFFUS												
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0005	æ	SFC.		4 0000 €			ש טטטטמה	EPISCH	0.02	900754	· F. • !	1 7000	100000	C.
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\$ 000	02200	7.6	0302	R 001460	-	2006	990000	40						1
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10100	ı	£.	First Section comput	COMPUTE	ES THE PRO	PROPER OF VALUE.	ALUE.							
50100	• ď	38	DAMERSION I (6)		161	161-prf16		***						1
100		SE .	- (Ca) -	TOTAL	TAPER I	THE CAR	1000	TO TOTAL TO THE CONTRACT OF TH	10104139					1
00100		3	COMMON/SING/LC. IR.NI	PINORIO	OT - NCS1 - NCS2	:52								
90100	.	83	COMPOST/BIZDCTT6.40.4	16.40.50	0) - SH(6-40)	3								1
		SE		72										
90111	::	33	DATA EPISLN/.0001/	7001										
00113	12:	25	SVSUHED										-	į
00114	13.	J=3	5											
51100	- 97	2	NELEN C. IN-0											
21100		3	SOMEOUT TO						30	17.55				

OI FOR								1			
	DEALC										
THIS COM	PILATION	AAS DONE	THIS COMPILATION AND DONE ON 19 DEC 66	-TV DATED 11 C 66 AT 21:06:12	^0^	1966 F4008					
SUBRO	SUBROUTINE ACALC		ENIRY POINT 000157	151000 1							
STORA	STORAGE USED ((JLOCK. 1	HAME LENGTH)	[H]							
	1	CODE 0	000225	•							
	0002	×	001676								
		٩	023160 00u001								
EXTER	NAL REFER	ENCES (9	EXTERNAL REFERENCES (BLUCK+ NAME)				Ø				
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STORAGE	16E ASSIGNAEUT	thent For	VARIABLES	(3L00Ke 1	151 AT 175	TYPE - MEI ATTUE : OCHTENA : MEI ATTUE					
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10.00		0									
10100		SUC.	1	. 1	KLA3+3+50+5T)						
10100		FIRS	FIRST SECTION COMPUTES	. •	ME PROPER OF VALUE	E (MEFUSION ESCATIONS ALDE	יוטני				
00103	;	UINE	UINCISTON ITES								
00104	ů	F.R.O.O	1014 YLD(5+1	COMMON YLD(5,40).: A:- E(6). MYOL	(6) . DFC (6)	**************************************	((c)+3(2)				
50100	• ~	はアピー	1) • VOC (40)	(30) • VOL (40) • T@• VOL (40) • TOTAN VOL (40) • TOTAN VOL (40) 1001 1001		•			1	:	
90100	9	FACU	05.78173CT	COMPONENT CONTRACTOR OF STREET			•				
00100	•	X > C U	T/dcoln1/nowhco	/L			Ţ.				
01100	10.	DIVE	DI VERSION DOLLY								
00111	11.	DATA	DATA EPISCN/.0001/	/1000							
00114	13.	SVSUMED	DUK	8)							
00115	14:	NEL EN	Z				İ				
00116	15.	INC'S									

00120	17.		N=:11+1(5)+1
00121	18.		IF (KLAG) 50, 50, 20
00124	19.	20	NIENZ+1
00125	20.		[F(N].61.41)60 to 60
00127	21. 20		NEP=:4EL+L-1
00130	22.		DO 10 KEN1521
00133	23.		SC#18C4+DFC (NFD*X) +1 (X)
45 TOO	24.		IF (ABS (SVSUM-SUM). LT. ASS (EPISLN-SVSUA)) 60 TO 13
00136	25.		XUSHKUS-2
00137	26.	9	CONTINUE
14100	27.	9	NEL=NEL-1
24100	26.		IF (NEL-1)11,12,13
00145	29.	13	N1=41+1
99100	20.		I+IC) I+INETE
24100	31.		1+747
00120	32.		GO TO 20
00151	33.	12	N1=41+1
25100	3		11216
00153	35.		6.3 To 20
95100	36.	÷	DIESUA
00155	37.		GHEOS
99100	38.		N-CONT-I(I)+I
00157	39.		DO 30 K=1.IR
29100	3		03(4)=1.0
00163	*DIAGNOSTIC*	•	THE TEST FOR EQUALITY RETREEN HOUSTNESSERS WAY NOT AS MARKETHER IN
00163	-14		IF (DY-6E-(1-6-9(K)-+2)) 60 19 30
99100	* 5.		DD(K)=DIFFUS(DT+R(K))
99100	43.	30	SUAHSUA+UD(K) +11(K)
00170	*		まついまの
1/100	45.		KETURN
00172	100		

SUBROUTINE DELTA ENTRY POINT 000172 SUBROUTINE DELTA ENTRY POINT 000172 STORAGE USED (JACK, NAME, LENGTM) 0001	DEC 66 AT 21:06:15 INT 000172 MGTH) LES THLOCK: TYPE: RELATIVE LOCATION: NEVE) 1 000152 17L 0003 R 00000 DCT 10 1 000152 17L 0003 R 00010 DCT 10 1 000154 NJ 0000 I 00015 ND4 10 1 00014 NJ 0000 I 00015 ND4 (6) DELTATIVE: RELATIVE LOCATION: NEVE) (6) DELTATIVE: NEVEL INCOME. NEVEL INCOME. NEVEL INCOME. NEVEL INCOME. NO. 1 0000 I 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NEVEL INCOME. NO. 1 00001 I 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NEVEL INCOME. NO. 1 00001 I 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NEVEL INCOME. NO. 1 00001 I 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NO. 1 000015 ND4 (6) DELTATIVE: NEVEL INCOME. NO. 1 000015 ND4 (6) DELTATIVE: NO. 1 000015 ND4 (6) DELTATI	0000 1 00000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 00000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 00000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 00000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000 1 0000	0.00 1
3ROUTINE DELTA 0001 •CODE 0000 •DATA 0002 •BLAN 0005 B1 0005 B1 0005 B1 0001 00056 153 000 1 000131 ISU 000 1 000131 ISU 2. 2. 2. 3. 4. 4. 4. 4. 4. 4. 4. 4. 4. 4. 4. 4. 4.	11106 1170 NO3	2 000000 1 000000 1 000000	616127 5 00:0016 7 1-00:0016 10:0010
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	.0ELT(6)		
5. 10. 11. 12. 12. 12. 13.			
11 10 10 10 10 10 10 10 10 10 10 10 10 1			
10. 11 10. 11 11. 14 12. 12			
10. 11. 14. 12. 12. 12. 12. 12. 12. 12. 12. 12. 12			
11. 14			
15. 12			
, ,			
15.			
00130 16. OLT=1.			
18.			
00134 19. 10 DELTCJIEDET4(1:=DEYCKDI) 00136 20. ISUMENT-1(1)	-DCY(NOI-NDK+NDK))		
22. 15			
.52			
25.	10 17		
00150 26. DELT(N)=DCY(L.ISUN.IT-1)	SUM-11-1		

28.	SUBROUTINE DELIA
00155 29. 1 00157 30. 00160 31.	DO 16 κ=1.η 16 Pλ=Pλ+OELT(Κ) κΕΤUΚι ΕΝΟ
END OF LISTING.	NG. 0 +JIAGNOSTIC+ AESSAGE(S).

							1	
COMPILAT THIS COM	OI FOR DIFFUS COMPILATION BY UNIVAC 1108 FORTRAH-IV SATES THIS COMPILATION 4AS DONE ON 19 DEC 66 AT 2	S NIVAC 110 AAS DONE	FORTRAH ON 19 DE		11 NOV 1966 F4005	6 F4009		
FUNC	FUNCTION DIFFUS		ENTRY POINT 000000	T 000000				
STOR	STORAGE USED (BLOCK.		NAME. LENGTH)	THI				
	0000 0000 0002	CUDE DO	000021 000021 000000					
EXTE	EXTERNAL REFERENCES (BLUCK+ NAME)	RENCES (3)	OCK - NAM	(3				
	0000	Surt						
STORAGE	IAGE ASSIGNAENT		FOR VARIABLES	(3LOCK. T	. RELATIVE	IPE. RELATIVE LOCATION. NEVEN		
0000	œ	000021 10L 000000 DIFFUS	0001 0004 R	000020 12L R 000050 EXP	0000	0000 9 000001 COPST	0000 8 000002 C1	3030 2 0C0003 C2
10100	:	FUNC	FUNCTION DIFFUSIOTER)	US(OT,R)				
00103	•	DATA X=0/	DATA COUSTICE	MARK COUSTICE/1:128379:	1. 50792710 - 9. H5968447	9.85960447		
01100		IFCX	IF (X-3-14) 10-10-20	02.50				
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13 ABSTRACT

A program is presented to compute fallout fission product absorption in an expanding, cooling, uniform field of gas and fallout particles where the rate of fission product absorption is controlled by fission product surface concentrations as given by Henry's law constants and diffusion of these fission products into the fallout particles. The calculations are made a nuclide chain at a time, employing nuclear device and fission product parameters. Program output includes average concentration versus particle size for each absorbed fission product at a preselected lower temperature and the amount of that fission product in the gas phase, with an option for the calculation of fission product radial profiles versus particle size.

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